# The $\ell\text{-adic}$ Fourier Transform

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## 1 Kloosterman sums

Let  $\psi \colon \mathbf{F}_p \to \mathbf{C}^{\times}$  or  $\overline{\mathbf{Q}}_p^{\times}$  be a character, e.g.

$$x \mapsto e^{2\pi i x/p}.$$

An important property is that

$$\sum_{a \in \mathbf{F}_p} \psi(ay) = \begin{cases} 0 & y \neq 0, \\ p & y = 0 \end{cases}$$
(1.1)

For  $a \in \mathbf{F}_p^*$ , the *Kloosterman sum* K(a) is

$$K(a) = \sum_{xy=a} \psi(x+y) = \sum_{x \in \mathbf{F}_p, x \neq 0} \psi(ax+x^{-1}).$$
(1.2)

**Example 1.1.** K(a) is always a real number, because the sum is symmetric with respect to complex conjugation. For p = 7 and a = 1, it is

$$\zeta_7^2 + \zeta_7^{-1}s + \zeta_7 + \zeta_7^{-2} + \zeta_7 + \zeta_7^{-1}$$

**Remark 1.2.** The analogue of K(a) over **R** would be something like

$$\int_{\mathbf{R}} e^{i(ax+1/x)} \, dx.$$

This isn't convergent, but

$$\int_{\mathbf{R}} e^{-(ax+1/x)} \, dx \sim \sqrt{a} K(\sqrt{a})$$

where K is a Bessel function. There is a parallel between these special functions and the special sheaves that will arise later.

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This was first studied by Kloosterman, in analyzing the Hardy-Littlewood circle method. Obviously |K(a)| < p; Kloosterman wanted to show that |K(a)| is much less than p. He showed that

$$|K(a)| \le p^{3/4}.$$

Weil later improved this to  $|K(a)| \leq 2\sqrt{p}$ .

We'll discuss Kloosterman's proof. He studied the sum

$$\sum_{a} |K(a)|^4$$

and showed

$$\sum_{a} |K(a)|^4 \le cp^3.$$

We start off by writing

$$\sum_{a} |K(a)|^4 = \sum_{a} \sum_{x,y,z,w} \psi(a(x+y-z-w) + (x^{-1}+y^{-1}-z^{-1}-w^{-1}))$$

Using (1.1), this is

$$= p \sum_{\substack{x+y=z+w \\ x^{-1}+y^{-1}\neq z^{-1}+w^{-1}}} \psi(x^{-1}+y^{-1}-z^{-1}-w^{-1}) + \# \left\{ \sum_{\substack{x+y=z+w \\ x^{-1}+y^{-1}\neq z^{-1}+w^{-1}}} \psi(x^{-1}+y^{-1}-z^{-1}-w^{-1}) + \# \left\{ \sum_{\substack{x-1+y-1=z^{-1}+w^{-1} \\ x^{-1}+y^{-1}\neq z^{-1}+w^{-1}}} \right\} \right)$$

Recall that we want to get a bound of  $p^3$ , while the trivial bound is about  $p^4$ . The second term  $\# \left\{ \substack{x+y=z+w\\ x^{-1}+y^{-1}=z^{-1}+w^{-1} } \right\}$  has only about  $p^2$  terms, so that's good. The first term is p times a sum over  $p^3$  things, so we need to do something intelligent there. Luckily, it has a *scaling* symmetry. The sum over each  $\mathbf{F}_p^*$ -orbits is -1 by (1.1).

So, letting 
$$N = \#\{x + y = z + w\}$$
 and  $A = \#\left\{ \substack{x+y=z+w\\x^{-1}+y^{-1}=z^{-1}+w^{-1}} \right\}$ , we get
$$\sum_{a} |K(a)|^4 = p\left(\frac{N-A}{p-1}(-1)+A\right).$$

To conclude that, note that N has size about  $p^3$ , A has size about  $p^2$ .

**Remark 1.3.** You can evaluate N and A exactly. For odd p, we think A = 3(p - 2)(p - 1). This gives

$$\sum_{a} |K(a)|^4 = 2p^3 + \text{ (lower order terms)}.$$

Kloosterman's argument is the earliest instance I know of the following principle:to bound a single value of a function, put that function in a family and raise it to a higher power. This idea was used again by Rankin in the context of modular forms, which Deligne said was an inspiration for his proof of the Weil conjectures.

## 2 The sheaf-function correspondence

We now want to implement this idea with sheaves. Let X be a variety over  $\mathbf{F}_p$ . Suppose you have a Weil sheaf  $\mathcal{F}$  on X, meaning a sheaf on  $X_{\overline{\mathbf{F}}_p}$  with a Frobenius endomorphism. Then we get a function f on |X| or  $X(\mathbf{F}_{p^n})$ , given by f(x) = trace of geometric Frobenius at x.

**Remark 2.1.** If for example  $\mathcal{F}$  is lisse and semisimple, then the function f determines the sheaf, because the Frobenii are dense in the monodromy group.

#### 2.1 Translation of sheaf-theoretic operations

Operations on sheaves can be translated into operations on functions.

• The tensor product of sheaves translates into product of functions.

$$\mathcal{F} \otimes \mathcal{G} \mapsto f_{\mathcal{F}} \cdot f_{\mathcal{G}}$$

• The pulblack of sheaves translates into pullback of functions.

$$\pi^*\mathcal{F} \mapsto f_{\mathcal{F}} \circ \pi$$

- If  $\mathcal{F}$  is pure of weight  $w \in \mathbb{Z}$ , then  $\mathcal{F}^{\vee}$  corresponds to the function  $\overline{f}p^{-w \cdot \deg}$ . For instance, the Kloosterman sums being real-valued corresponds to the Kloosterman sheaves being self-dual up to Tate twist.
- If  $\mathcal{F}$  is in the derived category of Weil sheaves<sup>1</sup> then

$$R\pi_{!}\mathcal{F}\mapsto \sum_{i}(-1)^{i}f_{H^{i}\mathcal{F}}$$

With these conventions, the Lefschetz trace formula translates into the the statement that the derived pushforward of sheaf corresponds to the pushforward of f as defined by

$$y \in Y(\mathbf{F}_{p^n}) \mapsto \sum_{x \in X(\mathbf{F}_{p^m}), \pi(x)=y} f(x).$$

<sup>&</sup>lt;sup>1</sup>Although we have glided over this point in this seminar, the construction of the "derived category of  $\ell$ -adic sheaves" (or Weil sheaves) is actually quite subtle. It is not obtained by the naïve construction taking the derived category of a category of  $\ell$ -adic sheaves, although this is often what one pretends for practical purposes. Suffice it to say that working rigorously with the "derived category of  $\ell$ -adic sheaves" requires a good deal more care than one might think; "arguments" which treat this category as a genuine derived category are merely reasoning by analogy.

For example, given a map  $\pi: X \to \mathbf{A}^1$ , then the function

$$y \in \mathbf{F}_{p^m} \mapsto \# X_y(\mathbf{F}_{p^m})$$

comes from  $R\pi_! \mathbf{Q}_{\ell}$ .

**Example 2.2.** We're going to make a sheaf corresponding to the function  $\psi$ .

We start out with the Artin-Schreier cover

$$y^p - y = x \subset \mathbf{A}^2.$$

This maps via the x-coordinate to  $\mathbf{A}^1$ , which is an étale cover. The Galois group is canonically  $\mathbf{Z}/p\mathbf{Z}$ , generated by  $y \mapsto y+1$ . In other words, this cover induces a map

$$\pi_1(\mathbf{A}^1_{\mathbf{F}_p}) \to \mathbf{Z}/p\mathbf{Z} \xrightarrow{\psi} \overline{\mathbf{Q}}_\ell^*.$$

Let  $\mathcal{L}_{\psi}$  be the corresponding rank 1 lisse sheaf on  $\mathbf{A}^1$ . We compute the associated function. We need to figure out where Frobenius goes. The Frobenius at  $x \in \mathbf{A}^1(\mathbf{F}_{p^m})$  takes  $(x, y) \mapsto (x^{p^m}, y^{p^m})$ . Of course  $x^{p^m} = x$ . We have

$$y^{p} = y + x$$
  

$$y^{p^{2}} = y^{p} + x^{p} = y + x + x^{p}$$
  

$$\vdots$$
  

$$y^{p^{m}} = y + x + x^{p} + \ldots + x^{p^{m-1}}$$

Therefore, the Frobenius at x takes

$$(x,y) \mapsto (x,y + \underbrace{x + x^p + x^{p^2} + \ldots + x^{p^{m-1}}}_{\operatorname{Tr}_{\mathbf{F}_p(x)}}).$$

The conclusion is that Frobenius acts on the stalk  $\mathcal{L}_{\psi}$  as multiplication by  $\psi(\operatorname{Tr}_{\mathbf{F}_{p^m}/\mathbf{F}_p}(x))$ . Therefore *geometric* Frobenius acts as multiplication by  $\psi(-\operatorname{Tr}_{\mathbf{F}_{p^m}/\mathbf{F}_p}(x)) = \overline{\psi}(\operatorname{Tr}_{\mathbf{F}_{p^m}/\mathbf{F}_p}(x))$ .

### 2.2 The method of families

Suppose we have an open subset  $U \subset \mathbf{A}^1$ , and  $\mathcal{G}$  is a Weil sheaf on U, associated to a function g. Assume  $g \geq 0$ . (This can be arranged by taking the sum of  $\mathcal{G}$  with its conjugate.)

We can then bound a single value of g by a sum:

$$g(x) \le \sum_{x \in \mathbf{F}_{p^n}} g(x).$$

Of course this isn't sharp, but if you apply this to large powers of g then it will be sharp.

To simplify things, assume  $H_c^0(G) = 0$ . Then

$$\sum_{x \in U(\mathbf{F}_{p^m})} g(x) = \sum_{\beta = \text{ eig. of } F \text{ on } H_c^2} \beta^m - \sum_{\alpha = \text{ eig. of } F \text{ on } H_c^1} \alpha^m.$$

We can easily analyze the  $\beta$ 's, by using Poincaré duality to relate  $H_c^2$  to the coinvariants of geometric  $\pi_1$ . But the point is that  $\max |\alpha| \leq \max |\beta|$ , which allows you to ignore  $\alpha$ . Why? This is because the expression is positive.

**Remark 2.3.** This same observation appears elsewhere. For instance, the Weil bound is not optimal. For a curve it gives  $p + 1 + 2g\sqrt{p}$ , but this can't be attained because it would give a negative number of points over  $\mathbf{F}_{p^2}$ .

This argument (applied to a high power of g) is the engine that provides the bounds in the proof.

The central part of the proof will be the following statement: if  $\mathcal{F}$  is pure of weight w on  $U \subset \mathbf{A}^1$ , then the weights of  $H^1_c(U, \mathcal{F})$  are  $\leq w + 1$ . In terms of the functions  $f_{\mathcal{F}}$  associated to  $\mathcal{F}$ , this statement translates to the bound

$$\sum_{x \in U(\mathbf{F}_{p^n})} f_{\mathcal{F}}(x) \le p^{m(w+1)/2}.$$

If  $H_c^0 = 0$  and  $H_c^2 = 0$  (it is easy to reduce to this case), then

$$\operatorname{Tr}(\operatorname{Frob}^m, H_c^1) = -\sum_{x \in U(\mathbf{F}_{p^n})} f_{\mathcal{F}}(x)$$

so this becomes a question of bounding the eigenvalues of Frobenius on cohomology.

To obtain this estimate, we try to embed it into a family. We want to find a function g on  $\mathbf{A}^1$  (associated to a sheaf) such that  $g(0 \in \mathbf{F}_{p^m}) = \sum f_{\mathcal{F}}(x)$ . Then we'll use the method of families. The punchline is that we take g to be the Fourier transform of f. Si next we'll make a sheaf associated to the function  $g = \mathrm{FT}(f)$ , defined by

$$\operatorname{FT}(f_{\mathcal{F}})(y) = \sum_{x \in \mathbf{F}_{p^m}} f_{\mathcal{F}}(x)\psi \circ \operatorname{Tr}(yx)$$

for  $y \in \mathbf{F}_{p^m}$ .

## 3 Fourier transform

Let  $\mathcal{F}$  be a sheaf on  $\mathbf{A}^1$ . We make a new sheaf  $\mathrm{FT}_{\psi} \mathcal{F}$  such that

$$f_{\mathrm{FT}_{\psi}}(y) = \sum_{x \in \mathbf{F}_{p^m}} f_{\mathcal{F}}(x)\psi \circ \mathrm{Tr}(yx).$$

We just replicate the Fourier transform step-by-step.

We start with  $\mathcal{F}$ , pull it back to  $\mathbf{A}^2$  via  $(x, y) \mapsto x$ . Then we tensor with  $m^* \mathcal{L}_{\psi}$ , where m(x, y) = xy. Finally, to sum over the first variable we push forward via  $(x, y) \mapsto y$ . The last step is to shift by degree 1, basically to preserve the property of being a sheaf (but it still might not quite).

Denote this functor by  $FT_{\psi}$ .

Theorem 3.1. We have

$$\operatorname{FT}_{\overline{\psi}} \circ FT_{\psi} = \operatorname{Id}(up \ to \ Tate \ twist).$$

This mirrors the usual calculation

$$\sum_{y} \psi(-yz) \sum_{x} f(x)\psi(yx) = \sum_{x,y} f(x)\psi(y(x-z))$$
$$= p^{m}f(z).$$

The proof replicates this calculation at the level of sheaves. The only step that wasn't formal was the calculation

$$\sum_{a \in \mathbf{F}_p} \psi(ay) = \begin{cases} 0 & y \neq 0\\ p & \end{cases}$$

so we need a sheaf-theoretic analogue of it, which is

$$H_c^*(\mathbf{A}_{\overline{\mathbf{F}}_p}^1, L_{\psi}) = 0.$$

To prove this, recall that the sheaf  $L_{\psi}$  came from the covering

$$y^p - y = x$$

by taking the  $\psi$ -component of the pushforward of the constant sheaf. Then  $H_c^*$  is the  $\psi^{\pm 1}$ -component of  $H_c^*(C, \overline{\mathbf{Q}}_{\ell})$ , which is 0 (since  $C = \mathbf{A}^1$ ).

The idea of the proof of the Weil conjectures is to bound F-eigenvalues on  $H^1_c(U \subset \mathbf{A}^1, \mathcal{G})$ , which is the fiber at 0 of  $\mathrm{FT}_{\psi}(\mathcal{G})$ .

## 4 Kloosterman sheaves

Recall that we defined the Kloosterman function

$$K(a) = \sum_{xy=a} \psi(x+y) = \sum_{x \in \mathbf{F}_p, x \neq 0} \psi(ax+x^{-1}).$$

We're going to make a sheaf Kl on  $\mathbf{G}_m$  such that

$$f_{\mathrm{Kl}}(a \in \mathbf{F}_{p^m}) = \sum_{x \in \mathbf{F}_{p^m}^*} \psi \circ \mathrm{Tr}(ax + x^{-1}).$$

We start with  $L_{\psi}$  on  $\mathbf{G}_m$  to get  $\psi(x)$ , apply inversion to get  $\psi(x^{-1})$ , and apply  $\mathrm{FT}_{\psi}$ .

This gives a lisse sheaf Kl on  $\mathbf{G}_m$ , pure of weight 1. Since rank(Kl) = 2, this corresponds to a representation  $\pi_1(\mathbf{G}_m) \to \operatorname{GL}_2(\overline{\mathbf{Q}}_\ell)$  whose Zariski closure is  $\operatorname{SL}_2$  (the real-ness suggests the sheaf is self-dual).

Suppose we want to understand Kl<sub>a</sub> for  $a \in \overline{\mathbf{F}}_p$ . Take a = 1. It is the  $\psi$ component of  $H_c^1(y^p - y = x + x^{-1})$ . We will show that dim  $H_c^1(y^p - y = x + x^{-1}) = 2(p-1) + 1$ . This strongly suggests that, because there are p-1 characters  $\psi$ , each
piece has dimension 2.

If you actually want to compute, you have to understand the behavior of the sheaf at  $\infty$ . Consider

$$y^p - y = x + x^{-1} \to \mathbf{G}_m$$

and compactify it to  $X \to \mathbf{P}^1$ , of degree p.

By Riemann-Hurwitz,

$$2g_X - 2 = p(-2) + \deg(\text{ram. divisor}).$$

The ramification is supported at  $0, \infty$ . Since the equation  $y^p - y = x + x^{-1}$  is symmetric, the answer will be the same at both points, so we just to the calculation 0. Localizing at 0, we need to consider the field extension L/K where  $K = \mathbf{F}_p((x))$ and L = K(y) with  $y^p - y = x + x^{-1}$ . Since v(x) = 1, we have v(y) = -1/p,  $\tau = y^{-1}$ is a uniformizer. The discriminant is the field extension

$$\prod_{i\neq j} (\tau_i - \tau_j)$$

Since the conjugates just add, a typical term is

$$\frac{1}{y+1} - \frac{1}{y} = \frac{-1}{y(y+1)}$$

with valuation 2/p. So the discriminant has valuation p(p-1)(2/p) = 2(p-1). (This is double what one would expect in characteristic 0.) So the conclusion is that

$$2g_X - 2 = p(-2) + 4(p-1) \implies \dim H^1(X) = 2(p-1).$$

To get C from the compactified guy, you delete 2 points so

$$\dim H^1(C) = 2(p-1) + 1.$$

So we've verified that  $\operatorname{rank} \mathrm{KL} = 2$ .

**Remark 4.1.** This is related to the fact that the K-Bessel function from Remark 1.2 satisfies a *second*-order differential equation.

Let's return to the estimate:

$$\sum |K(a)|^4 \sim 2p^3.$$

This sum can be interpreted as

$$\sum \operatorname{Tr}(F|H_c^i(\mathrm{Kl}\otimes \mathrm{Kl}^{\vee}\otimes \mathrm{Kl}\otimes \mathrm{Kl}^{\vee})(-1)^i.$$

By Deligne,  $H_c^1$  contributes a second-order term, so the leading term comes from

$$H_c^2 = (V \otimes V^{\vee} \otimes V \otimes V^{\vee})_{\pi_1^{\text{geom}}}(-1) = (V \otimes V^{\vee} \otimes V \otimes V^{\vee})^{\pi_1^{\text{geom}}}$$

We can interpret  $V \otimes V^{\vee} \otimes V \otimes V^{\vee} = \text{End}(V \otimes V^{\vee})$ . So we want to compute

dim End
$$(V \otimes V^{\vee})^{\pi_1^{\text{geom}}}$$
.

We have an irreducible decomposition of  $V \otimes V^{\vee}$  into the direct sum of a 3-dimensional representation and a 1-dimensional representation, so there are indeed two independent  $\pi_1^{\rm geom}$ -equivariant endomorphisms.