

MORAWETZ ESTIMATES FOR THE WAVE EQUATION AT LOW FREQUENCY

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ABSTRACT. We consider Morawetz estimates for weighted energy decay of solutions to the wave equation on scattering manifolds (i.e., those with large conic ends). We show that a Morawetz estimate persists for solutions that are localized at low frequencies, independent of the geometry of the compact part of the manifold. We further prove a new type of Morawetz estimate in this context, with both hypotheses and conclusion localized inside the forward light cone. This result allows us to gain a $1/2$ power of t decay relative to what would be dictated by energy estimates, in a small part of spacetime.

1. INTRODUCTION

In this paper, we show that the celebrated Morawetz estimate [14], expressing dispersion of solutions to the wave equation, holds for low-frequency solutions on a wide class of manifolds with asymptotically flat ends. We also generalize the estimate to a local-in-spacetime estimate inside the forward light cone.

It is well known that the decay of energy of a solution to $\square u = 0$ at *high frequencies* is closely tied to the geometry of geodesic rays; in particular, the existence of trapped geodesics is an obstruction to the uniform decay of local energy (see Ralston [18]). Many subsequent results have demonstrated that the decay of high frequency components of the solution persists in a wide variety of geometric settings in which there is no trapping of rays (see e.g. [15], [20], [21]).

On the other hand, the *low-frequency* behavior of the solutions is both more robust and less studied. Intuitively, long wavelengths should be sensitive only to the crudest aspects of the geometric setting, and in particular, trapping should not be an obstacle to decay. In this paper, we work in the setting of *scattering manifolds*, i.e., we assume that the manifold¹ X has ends resembling the large ends of cones, with the metric taking the form

$$(1.1) \quad g = dr^2 + r^2 h(r^{-1}, dy)$$

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¹We remark that in the body of the paper, X will refer to the manifold with boundary given by *compactifying* such a space.

where the noncompact ends are diffeomorphic to $(1, \infty)_r \times Y$, with Y a smooth manifold, and where h is a family (in r) of metrics on Y .² Our two main theorems demonstrate the insensitivity of the low-energy behavior of the wave equation to geometry on all but the largest scales. We also allow long-range perturbations of the metrics (1.1); see Section 2 for the precise definitions.

The results in this paper are as follows. To begin with, we state a result holding for low frequency solutions in spatial dimension $n \geq 4$ that generalizes the standard Morawetz estimate. We fix a cutoff function $\psi \in \mathcal{C}_c^\infty((0, \infty))$. Let Ψ_H denote the frequency-localization operator

$$\Psi_H = \psi(H^2(\Delta_g + V)).$$

We assume that the function r , appearing in the end structure of the metric (1.1), is extended to be a globally defined, smooth function on X , with $\min(r) \geq 2$ (so we may take $\log r$ with impunity). Let χ_0, χ_1 be a partition of unity with χ_1 supported in $r > R \gg 0$, hence in the ends in which the product decomposition (1.1) applies. Here and throughout the paper, $\|\bullet\|$ will refer to the norm in the space $L^2(X, dg)$. When we employ spacetime norms in §4 we will employ distinct notation $\|\bullet\|_M$ where M is (the compactification of) the spacetime $\mathbb{R} \times X$.

Theorem 1.1. *Let X be a scattering manifold of dimension $n \geq 4$ with a long-range metric, as in (2.1)–(2.2), and let $V \geq 0$ be a symbol of order $-2 - \nu$, $\nu > 0$. Let u be a solution to the inhomogeneous wave equation*

$$(\square + V)u = f, \quad f \in L^1(\mathbb{R}; L^2(X))$$

on $\mathbb{R}_t \times X$, with initial data

$$u|_{t=0} = u_0 \in H^1, \quad \partial_t u|_{t=0} = v_0 \in L^2.$$

Then for H sufficiently large,

$$\begin{aligned} (1.2) \quad & \int_0^\infty \left(\|r^{-3/2}\Psi_H u\|^2 + \|\chi_1(\log r)^{-1}r^{-1/2}\partial_r\Psi_H u\|^2 \right. \\ & \left. + \|\chi_1 r^{-3/2}\nabla_Y\Psi_H u\|^2 + \|\chi_0\nabla_g\Psi_H u\|^2 \right) dt \\ & \lesssim \mathcal{E}(0) + \left(\int_0^\infty \|f\| dt \right)^2, \end{aligned}$$

where $\mathcal{E}(0)$ is the initial energy

$$\mathcal{E}(0) = \|\Psi_H u_0\|_{H^1}^2 + \|\Psi_H v_0\|_{L^2}^2.$$

Here H^1 and L^2 denote the appropriately defined spaces on X (and $\|\bullet\|$ denotes the $L^2(X)$ norm): L^2 is the space of functions square integrable with respect to the metric, and H^1 , which will henceforth be written H_{sc}^1 (the “scattering” Sobolev space, associated to the geometry of large conic ends) denotes the space of functions satisfying

$$\int |u|^2 + |\nabla_g u|_g^2 dg < \infty.$$

²We may a priori assume that h is in fact a smooth tensor including dr components; that we may then change variables to remove these components is a result of Joshi-Sá Barreto [9], Section 2.

The Laplacian Δ_g is the nonnegative Laplacian, and $\psi(H^2\Delta_g)$ is localizing at low energies. The result thus shows that the dispersive effects of the large scale scattering geometry always hold for low-frequency solutions, regardless of trapping or other local features. We further remark that we could replace the $L^1(\mathbb{R}; L^2(X))$ norm on f on the RHS of the estimate by

$$(1.3) \quad \int_0^\infty \left\| (r^2 + t^2)^{1/4} \log(r^2 + t^2) f \right\|^2 dt,$$

if desired, to obtain a weighted L^2 spacetime norm instead.

Following the suggestion of an anonymous referee, we are in fact able to prove a refined version of this result, with logarithmic losses replaced by estimates in ℓ^1 and ℓ^∞ norms on energy in dyadic spatial shells. Let Υ_k denote a spatial decomposition in radial dyadic shells (to be discussed further in §3, see (3.18)), so $r \sim 2^k$ on Υ_k .

Theorem 1.1'. *Let X be a scattering manifold of dimension $n \geq 4$ with a long-range metric, as in (2.1)–(2.2), and let $V \geq 0$ be a symbol of order $-2 - \nu$, $\nu > 0$. Let u be a solution to the inhomogeneous wave equation*

$$(\square + V)u = f, \quad f \in L^1(\mathbb{R}; L^2(X)) \cap \ell^1(\mathbb{N}_k; L^2(\mathbb{R} \times \Upsilon_k))$$

on $\mathbb{R}_t \times X$, with initial data

$$u|_{t=0} = u_0 \in H^1, \quad \partial_t u|_{t=0} = v_0 \in L^2.$$

Then for H sufficiently large,

$$(1.4) \quad \begin{aligned} & \int_0^\infty \left(\left\| r^{-3/2} \Psi_H u \right\|^2 + \left\| \chi_1 r^{-3/2} \nabla_Y \Psi_H u \right\|^2 + \left\| \chi_0 \nabla_g \Psi_H u \right\|^2 \right) dt \\ & + \left\| \chi_1 r^{-1/2} \partial_r \Psi_H u \right\|_{\ell^\infty(\mathbb{N}_k; L^2([0, \infty)) \times \Upsilon_k)}^2 \\ & \lesssim \mathcal{E}(0) + \|f\|_{L^1([0, \infty); L^2(X))}^2 + \left\| r^{1/2} f \right\|_{\ell^1(\mathbb{N}_k; L^2([0, \infty) \times \Upsilon_k))}^2 \end{aligned}$$

where $\mathcal{E}(0)$ is the initial energy

$$\mathcal{E}(0) = \|\Psi_H u_0\|_{H^1}^2 + \|\Psi_H v_0\|_{L^2}^2.$$

Here we could drop $\int_0^\infty \left\| \chi_1 (\log r)^{-1} r^{-1/2} \partial_r \Psi_H u \right\|^2 dt$ from the left hand side of (1.4) as compared to (1.2) without any loss, since it can be estimated by

$$\begin{aligned} & \int_0^\infty \left\| \chi_1 (\log r)^{-1} r^{-1/2} \partial_r \Psi_H u \right\|^2 dt = \sum_k \int_0^\infty \int_{\Upsilon_k} \chi_1^2 (\log r)^{-2} r^{-1} |\partial_r \Psi_H u|^2 dg dt \\ & \lesssim \left(\sum_k k^{-2} \right) \left\| \chi_1 r^{-1/2} \partial_r \Psi_H u \right\|_{\ell^\infty(\mathbb{N}; L^2([0, \infty)) \times \Upsilon_k)}^2, \end{aligned}$$

so the quantity we estimate on the left hand side in Theorem 1.1' indeed stronger than that in Theorem 1.1. A similar argument shows that (1.3) also dominates the third term on the right hand side of (1.4).

Our second main result is a version of the Morawetz estimate that is localized in spacetime, with both hypotheses and conclusion localized in a sub-cone of the forward light cone. It holds in all dimensions $n \geq 3$.

For $\delta > 0$ let

$$(1.5) \quad \Omega = \Omega_\delta = \{t > 1/\delta, r/t < \delta\} \subset \mathbb{R} \times X.$$

This is the (asymptotic) cone in which we will localize.

In the statement of the following theorem, ∇ denotes the space-time gradient.

Theorem 1.2. *Let X be a scattering manifold of dimension $n \geq 3$ with a long-range metric as in (2.1)–(2.2), $V \geq 0$ a symbol of order $-2 - \nu$, $\nu > 0$. Let*

$$(\square + V)u = f.$$

Suppose that $\delta < 1$ and

$$\nabla \Psi_H u \in t^\kappa L^2(\Omega), \quad f \in r^{-1/2} t^{\kappa-1/2} L^2(\Omega), \quad \kappa \in \mathbb{R}.$$

Then for H sufficiently large,

$$\nabla \Psi_H u \in t^\kappa (r/t)^\sigma L^2(\Omega), \quad \Psi_H u \in t^\kappa r (r/t)^\sigma L^2(\Omega)$$

for $0 < \sigma < 1/2$.

In particular, if K is compact in X° , then on $\mathbb{R} \times K$, $\Psi_H u, \nabla \Psi_H u$ are in $t^{\kappa-1/2+\epsilon} L^2(\mathbb{R} \times X)$ for every $\epsilon > 0$ (for H sufficiently large).

This theorem gives almost half an order of decay as compared to the a priori assumption, but only in a small part of space-time. Indeed, for any $\delta_0 > 0$, the conclusion in $r/t > \delta_0$ is the same as hypothesis, so in particular δ in the statement of the theorem can be taken arbitrarily small without losing the strength of the conclusion. Thus, Theorem 1.2 really amounts to a regularity statement; it is a close analogue of the statement for hyperbolic PDE that over compact sets, one has extra regularity over a priori space-time Sobolev regularity (square integrability can be replaced by continuity along the flow). Indeed, an analogous statement is that locally L^2 solutions of the wave equation in space time, near $t = 0$, are in $t^\sigma L^2$ locally for $\sigma \in (0, 1/2)$; this follows from continuity in time with values in L^2 and the fact that $t^{-\sigma}$ is in $L^2(\mathbb{R})$ locally.

One virtue of this version of the Morawetz estimate, and even more of the approach to obtaining this estimate, is that they hold the promise of applying quite broadly, in non-product spacetimes. Even in the present inhomogeneous form they can be used to study the wave equation for metrics which arise by perturbing the Minkowski metric near infinity in the following way: pick p_j , $j = 1, \dots, N$, on the sphere at infinity, $\mathbb{S}^n = \partial \overline{\mathbb{R}_{t,z}^{n+1}}$, in the radial compactification of $\mathbb{R}^{n+1} = \mathbb{R}_t \times \mathbb{R}_z^n$ which lie in the interior of the forward light cone. Blow up the p_j ; if one of the p_j is the “north pole,” corresponding to $z = 0$, $t = +\infty$, a neighborhood of the front face can be identified with a neighborhood of “temporal infinity” in the product asymptotically Euclidean space-time (see Section 4); the other p_j can be transformed to this via isometries of Minkowski space-time. Thus, if the Minkowski metric is perturbed in a way corresponding to the local product structure these transformations induce, modulo decaying terms in time, the local Morawetz estimate is applicable by treating the decaying terms as inhomogeneities. In future work, we plan to address such non-product geometries more systematically.

By the standard energy estimate, if $(\square + V)u = 0$ and u is in the energy space, then for $\phi \in C_c^\infty(\mathbb{R})$,

$$\phi(r/t) \nabla u \in t^{1/2+\delta} L^2(\mathbb{R} \times X), \quad \delta > 0.$$

This observation immediately yields

Corollary 1.3. *Let X , g , and V be as in Theorem 1.2. Suppose that $\square u = 0$, with initial conditions*

$$u|_{t=0} = u_0 \in H^1, \quad \partial_t u|_{t=0} = v_0 \in L^2.$$

There exists $\delta > 0$ small such that

$$(1.6) \quad \iint_{\Omega_\delta} \left| t^{\sigma-1/2-\delta} r^{-1-\sigma} \Psi_H u \right|^2 + \left| t^{\sigma-1/2-\delta} r^{-\sigma} \nabla \Psi_H u \right|^2 dt dg \lesssim \mathcal{E}(0)$$

whenever $\sigma \in (0, 1/2)$, and $H \gg 0$ is sufficiently large.

The result is of special interest in the case $n = 3$, where Theorem 1.1 does not apply, but where Corollary 1.3 yields a spatially localized estimate. Let $K \subset X^\circ$ be compact. Then Corollary 1.3 specializes to show:

Corollary 1.4. *For all $\epsilon > 0$,*

$$\int_0^\infty \left\| t^{-\epsilon} \Psi_H u \right\|^2 + \left\| t^{-\epsilon} \nabla \Psi_H u \right\|^2 dt \lesssim \mathcal{E}(0).$$

for H sufficiently large.

Note, by contrast, that conservation of energy would give the same estimate with $t^{-\epsilon}$ replaced by $t^{-1/2-\epsilon}$, so this estimate represents an improvement of $t^{-1/2}$ in decay relative to energy estimates; by contrast, it loses $t^{-\epsilon}$ relative to Theorem 1.2 (which of course only holds for $n \geq 4$).

We emphasize that Theorem 1.2 and its corollaries can be proved without the need for localization at low frequencies so long as we can find a Morawetz commutant that works *globally*, and in particular this is the case on small perturbations of Euclidean space.

The proof of Theorem 1.1 is, as with the usual Morawetz inequality, a “multiplier” argument based on the first order differential operator

$$A_0 = (1/2)(\partial_r - \partial_r^*).$$

The subtlety is that the crucial commutator term

$$[\Delta, A_0]$$

is positive for $r \gg 0$, in the ends of the manifold, but certainly not in the interior where there is indeed no natural definition of the radial function r or of ∂_r . We must thus introduce cutoffs, which in turn introduce terms in the commutator that have to be treated as errors. At high frequency, these errors are disastrous, but at low frequency, we may employ a Poincaré/Hardy inequality to control them. Further technical subtleties arise in estimating remainder terms from the commutator, applied to $\Psi_H u$. These estimates are among the principal technical innovations here—see Proposition 2.1 below.

The proof of Theorem 1.2 uses a slightly different commutator argument (formally a “commutator” rather than a “multiplier” in the usual parlance) that occurs in space-time. The localizer in the cone Ω multiplies a Morawetz-like commutant with a positive weight in r but decay in t ; derivatives of the cutoff are controlled by the a priori decay assumptions, namely $\nabla \Psi_H u|_\Omega \in t^\kappa L^2(\mathbb{R} \times X)$, which in the corollaries are implied by energy estimates.

Related results have recently been pursued by a number of authors in the setting of asymptotically Euclidean manifolds. Bony-Häfner [1] explored Mourre estimates on asymptotically Euclidean spaces at low energy, while both Bouclet [3] and Bony-Häfner [2] have recently obtain iterated resolvent estimates in the low frequency regime sufficient to prove strong weighted decay estimates for asymptotically Euclidean metrics. These results yield stronger decay than what is obtained here, but in a narrower class of geometries and at the cost of a more intricate argument which

involves first proving resolvent estimates; one of the virtues of the techniques used here is the direct use of the hyperbolic equation, apart from the estimates necessary to localize at low energy. More in the spirit of our approach, Metcalfe-Tataru have proved an analogous theorem to our Theorem 1.1 for small perturbations of Euclidean space, and Tataru [19] has also proved strong energy decay estimates in a class of $3 + 1$ -dimensional Lorentz metrics that includes the Schwarzschild and Kerr spacetimes; see also the more recent [13]. (There has been considerable work on these specific Lorentzian examples motivated by problems in general relativity; we will not review that literature here.) The only prior results on low energy estimates in the setting of general scattering manifolds, in addition to the authors' [22], is the recent work of Guillarmou-Hassell-Sikora [6], which employs a sophisticated low-energy parametrix construction to obtain strong decay estimates. Due to the nature of the construction, the requirements on the metric in [6] are more stringent than the long-range assumptions used here.

Many of the low-frequency estimates used in this paper were first developed by the authors in [22] following the results of Bony-Häfner [1] in the asymptotically Euclidean setting, using rather different methods.

We have discussed both low- and high-frequency estimates above; for completeness, we remark that *intermediate frequencies* (i.e. frequencies in any compact subset of $(0, \infty)$) are always well-behaved in that the solution of the wave equation localized to these frequencies decays rapidly inside the forward light cone, in $r/t < c < 1$. This can be seen easily both by commutator methods (especially tools analogous to unreduced 2-body type problems, i.e. Mourre estimates as in the work of Mourre [16], Perry, Sigal and Simon [17]), Froese and Herbst [5], or directly using the resolvent of the Laplacian at the continuous spectrum, as described by Hassell and the first author in [7, 8] and the stationary phase lemma.

2. CONJUGATED SPECTRAL CUTOFFS

2.1. Notation and setting. Before stating our results on conjugated spectral cutoffs, we very briefly recall the basic definitions of the b- and scattering structures on a compact n -dimensional manifolds with boundary, denoted X ; we refer to [11] for more detail. A boundary defining function x on X is a non-negative C^∞ function on X whose zero set is exactly ∂X , and whose differential does not vanish there; a useful example to keep in mind is $x = r^{-1}$ (modified to be smooth at the origin) in the compactification of an asymptotically Euclidean space. We recall that $\dot{C}^\infty(X)$, which may also be called the set of Schwartz functions, is the subset of $C^\infty(X)$ consisting of functions vanishing at the boundary with all derivatives, the dual of $\dot{C}^\infty(X)$ is tempered distributional densities $\mathcal{C}^{-\infty}(X; \Omega X)$; tempered distributions $\mathcal{C}^{-\infty}(X)$ are elements of the dual of Schwartz densities, $\dot{C}^\infty(X; \Omega X)$.

Let $\mathcal{V}(X)$ be the Lie algebra of all C^∞ vector fields on X ; thus $\mathcal{V}(X)$ is the set of all C^∞ sections of TX . In local coordinates (x, y_1, \dots, y_{n-1}) ,

$$\partial_x, \partial_{y_1}, \dots, \partial_{y_{n-1}}$$

form a local basis for $\mathcal{V}(X)$, i.e. restrictions of elements of $\mathcal{V}(X)$ to the coordinate chart can be expressed uniquely as a linear combination of these vector fields with C^∞ coefficients. We next define $\mathcal{V}_b(X)$ to be the Lie algebra of C^∞ vector fields tangent to ∂X ; in local coordinates

$$x\partial_x, \partial_{y_1}, \dots, \partial_{y_{n-1}}$$

form a local basis in the same sense. Thus, $\mathcal{V}_b(X)$ is the set of all \mathcal{C}^∞ sections of a bundle, called the b-tangent bundle of X , denoted bTX . Finally, $\mathcal{V}_{sc}(X) = x\mathcal{V}_b(X)$ is the Lie algebra of scattering vector fields;

$$x^2\partial_x, x\partial_{y_1}, \dots, x\partial_{y_{n-1}}$$

form a local basis now. Again, $\mathcal{V}_{sc}(X)$ is the set of all \mathcal{C}^∞ sections of a bundle, called the scattering tangent bundle of X , denoted ${}^{sc}TX$. We write $\text{Diff}_b(X)$, resp. $\text{Diff}_{sc}(X)$, for the algebra of differential operators generated by $\mathcal{V}_b(X)$, resp. $\mathcal{V}_{sc}(X)$, over $\mathcal{C}^\infty(X)$; these are thus finite sums of finite products of vector fields in $\mathcal{V}_b(X)$, resp. $\mathcal{V}_{sc}(X)$, with $\mathcal{C}^\infty(X)$ coefficients. The dual bundles of bTX , resp. ${}^{sc}TX$ are ${}^bT^*X$ and ${}^{sc}T^*X$, and are called the b- and the scattering cotangent bundles, respectively; locally they are spanned by

$$\frac{dx}{x}, d_{y_1}, \dots, d_{y_{n-1}}, \text{ resp. } \frac{dx}{x^2}, \frac{d_{y_1}}{x}, \dots, \frac{d_{y_{n-1}}}{x}.$$

We let $S^k(X)$, the space of symbols of order k , consist of functions f such that

$$x^k Lf \in L^\infty(X) \text{ for all } L \in \text{Diff}_b(X).$$

We note, in particular, that

$$x^\rho \mathcal{C}^\infty(X) \subset S^{-\rho}(X)$$

since $\text{Diff}_b(X) \subset \text{Diff}(X)$. As $\text{Diff}_b(X)$ (a priori acting, say, on tempered distributions) preserves $S^k(X)$, and one can extend $\text{Diff}_b(X)$ and $\text{Diff}_{sc}(X)$ by ‘‘generalizing the coefficients:’’

$$S^k \text{Diff}_b^m(X) = \left\{ \sum_j a_j Q_j : a_j \in S^k(X), Q_j \in \text{Diff}_b^m(X) \right\},$$

with the sum being locally finite, and defining $S^k \text{Diff}_{sc}^m(X)$ similarly. In particular,

$$x^k \text{Diff}_b^m(X) \subset S^{-k} \text{Diff}_b^m(X), \quad x^k \text{Diff}_{sc}^m(X) \subset S^{-k} \text{Diff}_{sc}^m(X).$$

Then $Q \in S^k \text{Diff}_{sc}^m(X)$, $Q' \in S^{k'} \text{Diff}_{sc}^{m'}(X)$ gives $QQ' \in S^{k+k'} \text{Diff}_{sc}^{m+m'}(X)$, and the analogous statement for $S^k \text{Diff}_b^m(X)$ also holds.

In this paper, we are concerned with *scattering metrics*, that is to say, smooth metrics on X° that, near the boundary, take the form

$$(2.1) \quad g = \frac{dx^2}{x^4} + \frac{h}{x^2} + g_1$$

with $h = h(x, y, dy)$ a family of metrics on ∂X , and

$$(2.2) \quad g_1 \in S^{-\nu}(X; {}^{sc}T^*X \otimes {}^{sc}T^*X)$$

for some $\nu > 0$ is symmetric. Thus, g is a positive definite inner product on the fibers of ${}^{sc}TX$; long-range metrics like this were considered in [22]. We refer to Section 2 and the beginning of Section 3 of [22] for more details. Note that substituting $r = x^{-1}$ gives form 1.1 described above when $g_1 = 0$. We let $L^2(X)$ denote the space of functions square integrable with respect to the metric density; locally near a point in the boundary, this is equivalent to using the density

$$x^{-(n+1)} dx dy_1 \dots dy_{n-1}.$$

Note that the vector fields in $\mathcal{V}_{sc}(X)$ are precisely those with bounded length with respect to g . Correspondingly, a typical example of an element of $\text{Diff}_{sc}^2(X)$ is

the Laplace-Beltrami operator, $\Delta_g = d^*d$. We also permit short range potential perturbations of Δ_g , namely

$$(2.3) \quad V \in S^{-2-\nu}(X), \quad \nu > 0, \quad V \geq 0;$$

We assume that

$$V \geq 0,$$

though the arguments also work for potentials with small negative parts; see [22, Footnote 5].

2.2. Estimates. In this section we employ some of the results of [22] to obtain the following. Let $\psi \in \mathcal{C}_c^\infty(\mathbb{R})$ and

$$\Psi_H = \psi(H^2(\Delta_g + V)), \quad \Delta = \Delta_g.$$

Proposition 2.1. *Suppose $n \geq 3$, g and V as in (2.1), (2.2), (2.3). For $0 \leq s$, $0 \leq \rho$, $s + \rho < \min(2, n/2)$, $L \in \text{Diff}_{\text{sc}}^1(X)$,*

$$Lx^{s+\rho}\Psi_Hx^{-s} \leq CH^{-\rho}, \quad H > 1.$$

While the spectral cutoff is in the “(anti-)semi-classical” operator $H^2(\Delta_g + V)$ with large parameter, we emphasize that the operator L in the above proposition is simply a differential operator without parameter. As first-order scattering differential operators are spanned over $\mathcal{C}^\infty(X)$ by $x^2\partial_x$, ∂_y , and the constant function, to prove the proposition, it suffices to check for these particular values of L . Equivalently, we can simply use the vector-valued $L = \nabla_g$, the gradient with respect to the scattering metric, as well as $L = 1$.

Note that for $\rho = 0$ and $s = 0$ Proposition 2.1 follows automatically from the functional calculus and elliptic regularity; for $\rho = 0$ and $s = 1$ this is proved in [22], hence it follows by interpolation for $\rho = 0$, $0 \leq s \leq 1$. Thus, if $\rho = 0$, we need to deal with $1 < s < 2$ if $n \geq 4$, and $1 < s < 3/2$ if $n = 3$.

Now, with $\tilde{\psi} \in \mathcal{C}_c^\infty(\mathbb{C})$ an almost analytic extension of ψ , and

$$R(z) = (H^2(\Delta_g + V) - z)^{-1}$$

we have

$$(2.4) \quad Lx^{s+\rho}\Psi_Hx^{-s} = \frac{1}{2\pi} \int \bar{\partial}\tilde{\psi}(z)Lx^{s+\rho}R(z)x^{-s} dz d\bar{z}.$$

Thus, we only need to obtain uniform bounds on $Lx^{s+\rho}R(z)x^{-s}$ to prove the proposition.

We begin with a preliminary lemma, which is an analogue of [22, Proposition 4.3], where it was proved in the range $0 \leq s < 1/2$ for $n \geq 3$.

Lemma 2.2. *For $0 \leq s < \min(1, (n-2)/2)$,*

$$\|x^s\nabla_g u\| \leq C\|(\Delta_g + V)u\|^{(1+s)/2}\|u\|^{(1-s)/2}$$

Proof. By [22, Proposition 4.3], we may assume $n \geq 4$; then the range in the lemma is $0 \leq s < 1$. By [22, Equation (A.5)], we have for $0 \leq s < (n-2)/2$,

$$(2.5) \quad \|x^s\nabla_g u\| \leq C\|(\Delta_g + V)u\|^{1/2}\|x^{2s}u\|^{1/2}.$$

Moreover, by [22, Corollary 3.5], we have for $0 < s < (n-2)/2$, $0 \leq \theta \leq 1$,

$$(2.6) \quad \|x^{s+\theta}u\| \leq C\|x^s\nabla u\|^\theta\|x^s u\|^{1-\theta}.$$

Applying this with $0 < \theta = s < 1$, we obtain that

$$\|x^{2s}u\| \leq C\|x^s\nabla_g u\|^s\|x^s u\|^{1-s}.$$

Substituting into (2.5) yields

$$(2.7) \quad \|x^s\nabla_g u\|^{1-s/2} \leq C\|(\Delta_g + V)u\|^{1/2}\|x^s u\|^{(1-s)/2}.$$

Now using (2.6) with 0 in place of s and s in place of θ yields

$$\|x^s u\| \leq C\|\nabla_g u\|^s\|u\|^{1-s} \leq C\|(\Delta_g + V)u\|^{s/2}\|u\|^{1-s/2},$$

with the second inequality arising from

$$\|\nabla_g u\|^2 \leq \|(\Delta_g + V)u\| \|u\|,$$

which is (2.5) with $s = 0$. Substitution into (2.7) yields

$$(2.8) \quad \|x^s\nabla_g u\|^{1-s/2} \leq C\|(\Delta_g + V)u\|^{1/2+s(1-s)/4}\|u\|^{(1-s)(1-s/2)/2}.$$

Since $1/2+s(1-s)/4 = (1-s/2)(1+s)/2$, raising both sides to the power $(1-s/2)^{-1}$ yields

$$\|x^s\nabla_g u\| \leq C\|(\Delta_g + V)u\|^{(1+s)/2}\|u\|^{(1-s)/2}.$$

This finishes the proof. \square

Now, letting $u = R(z)f$, and writing

$$(\Delta_g + V)u = H^{-2}(H^2(\Delta_g + V) - z)u + H^{-2}zu,$$

we deduce from Lemma 2.2 that for $0 \leq s < \min(2, (n-2)/2)$,

$$\|x^s\nabla_g R(z)f\| \leq C\|H^{-2}(f + zR(z)f)\|^{(1+s)/2}\|R(z)f\|^{(1-s)/2}.$$

Since $\|R(z)f\| \leq |\operatorname{Im} z|^{-1}\|f\|$, this gives

$$(2.9) \quad \|x^s\nabla_g R(z)f\| \leq CH^{-(1+s)}(1 + |z|/|\operatorname{Im} z|)^{(1+s)/2}|\operatorname{Im} z|^{-(1-s)/2}\|f\|$$

$$(2.10) \quad \leq CH^{-(1+s)}|z|^{(1+s)/2}|\operatorname{Im} z|^{-1}\|f\|.$$

Using the Hardy/Poincaré inequality [22, Proposition 3.4], we deduce the following extension of [22, Proposition 4.5] (where it was proved for $0 \leq s < 1/2$, i.e. the result below is only an improvement if $n \geq 4$):

Corollary 2.3. *Suppose $0 \leq s < \min(1, (n-2)/2)$, $L \in S^{-1-s}\operatorname{Diff}_b^1(X)$. Then*

$$\|LR(z)f\| \leq CH^{-(1+s)}|z|^{(1+s)/2}|\operatorname{Im} z|^{-1}\|f\|.$$

We can extend the range of s when L is zero'th order by interpolating the case $s = 0$ above, namely

$$\|xR(z)f\| \leq CH^{-1}|z|^{1/2}|\operatorname{Im} z|^{-1}\|f\|,$$

with the estimate $\|R(z)f\| \leq |\operatorname{Im} z|^{-1}\|f\|$ to obtain the following result (in which the range $1 \leq s < \min(2, n/2)$ is a special case of Corollary 2.3, and it is the range $0 \leq s \leq 1$ that follows from the interpolation).

Corollary 2.4. *Suppose $0 \leq s < \min(2, n/2)$. Then*

$$\|x^s R(z)f\| \leq CH^{-s}|z|^{s/2}|\operatorname{Im} z|^{-1}\|f\|.$$

As we are interested in the functional calculus for compactly supported functions, we suppress the large z behavior, and work with z in a compact set. As a consequence of the preceding two corollaries we conclude:

Corollary 2.5. For $0 \leq s < \min(2, n/2)$, $L \in S^{-s}\text{Diff}_{\text{sc}}^1(X)$, z in a compact set,

$$(2.11) \quad \|LR(z)f\| \leq CH^{-s}|\text{Im } z|^{-1}\|f\|,$$

uniformly in z and $H \geq 1$.

Proof. If $1 \leq s$, then $L \in S^{-s}\text{Diff}_{\text{sc}}^1(X) \subset S^{-s}\text{Diff}_{\text{b}}^1(X)$ and Corollary 2.3 proves the result.

On the other hand, if $0 \leq s < 1$, one can write $L = aV + b$, $a, b \in S^{-s}(X)$, $V \in \mathcal{V}_{\text{sc}}(X) = x\mathcal{V}_{\text{b}}(X)$, so $aV \in S^{-1-s}\text{Diff}_{\text{b}}^1(X)$, so Corollary 2.3 gives

$$\|aVR(z)f\| \leq CH^{-1-s}|\text{Im } z|^{-1}\|f\|.$$

Furthermore, Corollary 2.4 yields

$$\|bR(z)f\| \leq CH^{-s}|\text{Im } z|^{-1}\|f\|.$$

Combining these two estimates proves the corollary. \square

Now we consider conjugates of $R(z)$. Below sometimes $n = 3$ and $n \geq 4$ are treated separately because of the stronger limitations $n = 3$ imposes in the preceding results: $\min(2, n/2) = 3/2$, rather than 2, in that case. First we prove a result for conjugation by small powers of x (namely, x^s , $0 \leq s \leq 1$); then in Proposition 2.7 we increase the range. Roughly speaking, the argument of Proposition 2.7 reduces the weight by 1 (and could be done inductively if Corollaries 2.3 and 2.5 had been proved for an extended range of weights); Lemma 2.6 deals with the fractional part of the weight (and thus would be the base case in the inductive argument).

Lemma 2.6. Suppose $n \geq 3$. Suppose $0 \leq s \leq 1$, $0 \leq \rho$, $s + \rho < \min(2, n/2)$, and if $n = 3$ then in addition $\rho < 1$. For $L \in \text{Diff}_{\text{sc}}^1(X)$, and z in a compact set,

$$(2.12) \quad \|Lx^{\rho+s}R(z)x^{-s}\| \leq CH^{-\rho}|\text{Im } z|^{-2}$$

uniformly as $H \rightarrow +\infty$.

Proof. The case $s = 0$ follows from Corollary 2.5. Note that the conditions imply $0 \leq \rho < 2$ if $n \geq 4$.

If $n = 3$, assume that $1/2 < s \leq 1$; if $n \geq 4$ no additional assumption is made. Then

$$\begin{aligned} Lx^{s+\rho}R(z)x^{-s} &= Lx^\rho R(z) + Lx^{s+\rho}[R(z), x^{-s}] \\ &= Lx^\rho R(z) - Lx^{s+\rho}R(z)[H^2(\Delta + V), x^{-s}]R(z). \end{aligned}$$

For $L \in \text{Diff}_{\text{sc}}^1(X)$, by Corollary 2.5, as $0 \leq \rho < \min(2, n/2)$,

$$\|Lx^\rho R(z)\| \leq CH^{-\rho}|\text{Im } z|^{-1}$$

As for the second term on the right hand side, we note that $[\Delta + V, x^{-s}] \in x^{2-s}\text{Diff}_{\text{b}}^1(X)$. Consequently, we may employ Corollary 2.3, using $0 < s \leq 1$ (so $1 \leq 2 - s < 2$) if $n \geq 4$, resp. $1/2 < s \leq 1$ (so $1 \leq 2 - s < 3/2$) if $n = 3$. This yields

$$\|[\Delta + V, x^{-s}]R(z)\| \leq CH^{-(2-s)}|\text{Im } z|^{-1}.$$

On the other hand, as $0 \leq s + \rho < \min(2, n/2)$, by Corollary 2.5,

$$\|Lx^{s+\rho}R(z)\| \leq CH^{-(s+\rho)}|\text{Im } z|^{-1},$$

so

$$\|Lx^{s+\rho}R(z)[H^2(\Delta + V), x^{-s}]R(z)\| \leq CH^{-\rho}|\text{Im } z|^{-2}.$$

Combining these results proves (2.12) if $0 < s \leq 1$ and $n \geq 4$, and also if $1/2 < s \leq 1$ and $n = 3$, and completes the proof in these cases.

It remains to deal with $n = 3$, and $0 < s \leq 1/2$. This follows by interpolation between $s = 0$ and $s > 1/2$, noting that given $0 \leq \rho < 1$ we can take $s' \in (1/2, 1)$ such that $\rho + s' < 3/2$; we then interpolate between 0 and this value s' . This last step is the reason for the restriction $\rho < 1$; if $\rho \geq 1$, the desired s' does not exist. This completes the proof of the Lemma. \square

We now extend the range of allowable exponents s by an inductive argument:

Proposition 2.7. *Suppose $n \geq 3$. Suppose $\rho \geq 0$, $L \in \text{Diff}_{\text{sc}}^1(X)$. For $0 \leq s \leq s + \rho < \min(2, n/2)$,*

$$(2.13) \quad \|Lx^{s+\rho}R(z)x^{-s}\| \leq CH^{-\rho}|\text{Im } z|^{-4}.$$

Proof. If $s = 0$, the result follows from Corollary 2.5. If $n \geq 4$, $0 < s < 1$, it follows from Lemma 2.6.

Assume now that $1 \leq s < \min(2, n/2)$. If $n = 3$, let $\rho' = s/2$; if $n \geq 4$, let $\rho' = s - 1$. In either case, $0 \leq \rho' < 1$; if $n = 3$ then in addition $\rho' > 1/2$. Let

$$L' = Lx\rho \in x^\rho \text{Diff}_{\text{sc}}^1(X).$$

Write

$$\begin{aligned} L'x^sR(z)x^{-s} &= L'R(z) + L'x^s[R(z), x^{-s}] \\ &= L'R(z) - \left(L'x^sR(z)x^{-\rho'}\right) \left(x^{\rho'}[H^2(\Delta + V), x^{-s}]R(z)\right). \end{aligned}$$

First, by Corollary 2.5 since $\rho' < \min(n/2, 2)$,

$$\|L'R(z)\| \leq CH^{-\rho'}|\text{Im } z|^{-2}.$$

Next, we apply (2.12), with $s = \rho'$ and ρ replaced by $\rho + s - \rho'$. The hypotheses are satisfied as $\rho' + (\rho + s - \rho') = \rho + s < \min(n/2, 2)$ and as $0 \leq \rho + s - \rho' < 1$ if $n = 3$ since $\rho + s < 3/2$, $\rho' = s/2$ and $s \geq 1$, resp. $0 \leq \rho + s - \rho' = \rho + 1$ if $n \geq 4$,

$$\|L'x^sR(z)x^{-\rho'}\| \leq CH^{-s-\rho+\rho'}|\text{Im } z|^{-2}.$$

On the other hand, $x^{\rho'}[\Delta + V, x^{-s}] \in x^{2-s+\rho'}\text{Diff}_{\text{b}}^1(X)$, so by Corollary 2.3 (taking into account that if $n \geq 4$, $2 - s + \rho' = 1$, and if $n = 3$ then $1 \leq 2 - s + \rho' < 3/2$),

$$\|x^{\rho'}[(\Delta + V), x^{-s}]R(z)\| \leq CH^{-2+s-\rho'}|\text{Im } z|^{-2}.$$

Combining these results, we have shown that

$$\|L'x^sR(z)x^{-s}\| \leq CH^{-\rho}|\text{Im } z|^{-4},$$

completing the proof if $1 \leq s < \min(2, n/2)$. If $n \geq 4$, we already covered the case of $s < 1$ in Lemma 2.6, hence if $n \geq 4$, the proof is complete.

So suppose that $n = 3$. If $0 < s < 1$, and $0 \leq \rho < 1/2$, then one can interpolate between $s = 0$ and $s = 1$ with the same (fixed) value of ρ to obtain (2.13), completing the proof if $0 \leq \rho < 1/2$ in the full range of s , $0 \leq s < 3/2$ (subject to $s + \rho' < 3/2$). Finally, with $0 \leq s + \rho' < 3/2$ fixed, one can interpolate between $s = 0$ and $\rho = 0$ to obtain the full result. \square

In view of (2.4), Proposition 2.7 proves Proposition 2.1.

We need a second, much less delicate, property of spectral cutoffs. Note that here we do not need to work with low energies: the cutoff ψ is fixed, and there is

no H in the statement of the proposition. (If we worked with $\psi(H^2(\Delta + V))$, H could be traced through the argument given below to yield polynomially growing bounds in H , i.e. one can gain decay in t at the cost of losing powers of H .)

Proposition 2.8. *Suppose that $\phi, \chi \in \mathcal{C}_c^\infty(\mathbb{R})$ and $\text{supp}(1 - \chi) \cap \text{supp } \phi = \emptyset$, $\psi \in \mathcal{C}_c^\infty(\mathbb{R})$. Also let $r = x^{-1}$. Then for $L \in \text{Diff}_{\text{sc}}^1(X)$, $N \in \mathbb{N}$,*

$$\|L\phi(r/t)\psi(\Delta + V)(1 - \chi(r/t))\|_{\mathcal{L}(L^2, L^2)} \leq C_N t^{-N}, \quad t \geq 1.$$

Proof. We first remark that

$$[x^2 \partial_x, \phi(1/(xt))] = -t^{-1} \phi'(1/(xt)),$$

while $x D_{y_j}$ and elements of $\mathcal{C}^\infty(X)$ commute with the multiplication operator by $\phi(1/(xt))$. Thus, by an inductive argument, for $Q \in \text{Diff}_{\text{sc}}^m(X)$, we have

$$(2.14) \quad [Q, \phi(1/(xt))] = \sum_{|\alpha|+j \leq m-1} \sum_k f_{j,\alpha,k}(t) \phi_{j,\alpha,k}(1/(xt)) a_{j,\alpha,k}(x^2 D_x)^j (x D_y)^\alpha$$

where the sum over k is finite, $f_{j,\alpha,k} \in S^{-1}([1, \infty))$ (i.e. is \mathcal{C}^∞ and is a symbol of order -1 at infinity), $a_{j,\alpha,k} \in \mathcal{C}^\infty(X)$, $\phi_{j,\alpha,k} \in \mathcal{C}_c^\infty(\mathbb{R})$ and $\text{supp } \phi_{j,\alpha,k} \subset \text{supp } d\phi$. Equivalently, we may put all of the factors $\phi_{j,\alpha,k}(1/(xt))$ on the right (at the cost of changing these factors as well as the other coefficients), so

$$(2.15) \quad [Q, \phi(1/(xt))] = \sum_{|\alpha|+j \leq m-1} \sum_k \tilde{f}_{j,\alpha,k}(t) \tilde{a}_{j,\alpha,k}(x^2 D_x)^j (x D_y)^\alpha \tilde{\phi}_{j,\alpha,k}(1/(xt)),$$

with the tilded objects having the same properties as the untilded ones above.

For each $t \geq 1$ we may write, $L\phi(r/t) = \phi(r/t)L + [L, \phi(r/t)]$ and use that for any $Q \in \text{Diff}_{\text{sc}}^m(X)$, $Q\psi(\Delta + V)$ is bounded on L^2 (by elliptic regularity), so an immediate consequence of (2.14) is that $L\phi(r/t)\psi(\Delta + V)(1 - \chi(r/t))$ is a bounded operator on L^2 which is uniformly bounded in $t \geq 1$; we now must obtain improved (decaying) bounds in t .

To do so, we work with the resolvent $R(z) = (\Delta + V - z)^{-1}$, $\text{Im } z \neq 0$, and note that if $\rho \in \mathcal{C}_c^\infty(\mathbb{R})$ and $1 - \chi$ have disjoint support, so $\rho(1 - \chi) = 0$, then commuting ρ through the resolvent,

$$(2.16) \quad \begin{aligned} \rho(1/(xt))R(z)(1 - \chi(1/(xt))) &= -[R(z), \rho(1/(xt))](1 - \chi(1/(xt))) \\ &= R(z)[\Delta + V, \rho(1/(xt))]R(z)(1 - \chi(1/(xt))). \end{aligned}$$

By (2.15), this has the form

$$(2.17) \quad \sum_{|\alpha|+j \leq 1} \sum_k f_{j,\alpha,k}(t) R(z) a_{j,\alpha,k}(x^2 D_x)^j (x D_y)^\alpha \rho_{j,\alpha,k}(1/(xt)) R(z)(1 - \chi(1/(xt))),$$

with $\text{supp } \rho_{j,\alpha,k} \subset \text{supp } d\rho$. Thus, $\rho_{j,\alpha,k}(1/(xt))R(z)(1 - \chi(1/(xt)))$ is of the same form as the left hand side of (2.16), so (2.16) can be further expanded by expanding these last three factors in (2.17) as in (2.16). Note that as we expand, each time we obtain a factor like $f_{j,\alpha,k} \in S^{-1}([1, \infty))$; these commute with all other factors, and after N iterations, the product of these is in $S^{-N}([1, \infty))$. Inductively, doing N iterations, we deduce that $\phi(1/(xt))R(z)(1 - \chi(1/(xt)))$ is a sum of terms of the form

$$(2.18) \quad f(t)R(z)A_1 R(z)A_2 \dots R(z)A_N \tilde{\phi}(1/(xt))R(z)(1 - \chi(1/(xt))),$$

with $A_j \in \text{Diff}_{\text{sc}}^1(X)$, $\tilde{\phi} \in \mathcal{C}_c^\infty(\mathbb{R})$ with $\text{supp } \tilde{\phi} \subset \text{supp } d\phi$, and $f \in S^{-N}([1, \infty))$, and one can also interchange A_N and $\tilde{\phi}(1/(xt))$ if convenient (at the cost of obtaining different operators in the same class).

But $\|A_j R(z)\| \leq C |\text{Im } z|^{-1}$ for z in a compact set, and similarly for $LR(z)$ when $L \in \text{Diff}_{\text{sc}}^1(X)$ (see Corollary 2.5 with H fixed) so using that $\tilde{\phi}$ and $1 - \chi$ are uniformly bounded in sup norm, hence as bounded operators on L^2 , we deduce that for $L \in \text{Diff}_{\text{sc}}^1(X)$,

$$\|L\phi(1/(xt))R(z)(1 - \chi(1/(xt)))\| \leq Ct^{-N} |\text{Im } z|^{-N-1}.$$

The Cauchy-Stokes formula, with $\tilde{\psi} \in \mathcal{C}_c^\infty(\mathbb{C})$ an almost analytic extension of ψ ,

$$(2.19) \quad \begin{aligned} & L\phi(1/(xt))\psi(\Delta + V)(1 - \chi(1/(xt))) \\ &= \frac{1}{2\pi} \int \bar{\partial} \tilde{\psi}(z) L\phi(1/(xt))R(z)(1 - \chi(1/(xt))) dz d\bar{z}, \end{aligned}$$

now immediately proves the proposition. \square

In fact, it is useful to add a weight in x as well:

Proposition 2.9. *Suppose that $\phi, \chi \in \mathcal{C}_c^\infty(\mathbb{R})$ and $\text{supp}(1 - \chi) \cap \text{supp } \phi = \emptyset$, $\psi \in \mathcal{C}_c^\infty(\mathbb{R})$. Also let $r = x^{-1}$. Then for $L \in \text{Diff}_{\text{sc}}^1(X)$, $N \in \mathbb{N}$, $m \in \mathbb{N}$*

$$\|Lr^m \phi(r/t)\psi(\Delta + V)(1 - \chi(r/t))r^m\|_{\mathcal{L}(L^2, L^2)} \leq C_N t^{-N}, \quad t \geq 1.$$

Proof. First, the $r^m = x^{-m}$ in front of $\psi(\Delta + V)$ is harmless since on the support of ϕ it is bounded by t^m , so it can be absorbed into t^{-N} on the right hand side if we increase N . Moreover, for each t , $Qx^m \psi(\Delta + V)x^{-m}$ is bounded for $Q \in \text{Diff}_{\text{sc}}^1(X)$, as is immediate from the scattering calculus of Melrose as $\psi(\Delta + V)$ is a pseudodifferential operator of order $-\infty$. However, there is a simple direct argument: via the Cauchy-Stokes formula this reduces to a boundedness statement for $Qx^m R(z)x^{-m}$. This in turn is shown by commuting x^{-m} through $R(z)$, much as we commuted $\phi(r/t)$ above, in this case gaining a power of x (instead of t^{-1}) each time we do a commutation, so in m steps we reduce to a product of $m + 1$ resolvents and bounded functions, giving a bound

$$(2.20) \quad \|Qx^m R(z)x^{-m}\| \leq C |\text{Im } z|^{-m-1}.$$

Thus, as in the previous proposition, we only need to prove the decaying uniform estimate in t . (Strictly speaking, one has to regularize, replacing x^{-m} by $(1 + \epsilon x^{-1})^{-m} x^{-m}$, which is bounded for $\epsilon > 0$, and let $\epsilon \rightarrow 0$, in which case the commutator argument presented above gives uniform control of the terms.)

We thus proceed as above, so by (2.18) we only need to prove that

$$A_N \tilde{\phi}(1/(xt))R(z)(1 - \chi(1/(xt)))x^{-m}$$

is L^2 bounded. As in the fixed t setting, it is convenient to insert a factor of x^m on the left, using that $x^{-m} \lesssim t^m$ on the support of $\tilde{\phi}$, and thus, commuting $\tilde{\phi}$ through A_N , and using the uniform boundedness of $\tilde{\phi}$ and $1 - \chi$ in the sup norm, we are reduced to showing that $A'_N x^m R(z)x^{-m}$ is bounded (with a polynomial bound in $|\text{Im } z|^{-1}$). But this is exactly the content of (2.20), completing the proof. \square

3. GLOBAL MORAWETZ ESTIMATE

In this section, we return to the notation $r = 1/x$ to facilitate comparisons with the usual method of Morawetz estimates. To simplify the computation, we assume $g_1 = 0$ with the notation of (2.1)–(2.2). Then, in the paragraph of (3.7) we reinstate arbitrary g_1 (with $\nu > 0$). Thus in the present notation, we have

$$(3.1) \quad g = g_0 = dr^2 + r^2 h(1/r, y, dy)$$

with $r \in (r_0, \infty)$. Then we compute

$$(3.2) \quad \begin{aligned} \partial_r^* &= -\partial_r - \partial_r(\log(r^{n-1}\sqrt{h})) = -\partial_r - \frac{n-1}{r} - \frac{1}{2}\partial_r \log h(1/r) \\ &= -\partial_r - \frac{n-1}{r} + e \in r^{-1}\text{Diff}_b^1(X) \end{aligned}$$

with $e \in S^{-2}(X)$.

We as usual may write, in the ends of X where $r \gg 0$,

$$\Delta_g = \partial_r^* \partial_r + r^{-2} \Delta_Y \in r^{-2} \text{Diff}_b^2(X)$$

where Δ_Y is the family of Laplacians on the boundary at infinity defined by the family of metrics $h(1/r, \bullet)$.

Now in general, for $F(r) \in \mathcal{C}^\infty([0, \infty))$ we let

$$(3.3) \quad A_F = (1/2)(F(r)\partial_r - \partial_r^* F(r)),$$

which is thus skew-adjoint. The principal symbol of the commutator of Δ_g with A_F is straightforward to compute, and one deduces that $[\Delta, A_F]$ and $2\partial_r^* F'(r)\partial_r + 2r^{-3}F(r)\Delta_Y$ have the same principal symbol, so their difference is first order. Since they are both formally self-adjoint and real, the same holds for the difference, which is thus zero'th order, and can be computed by applying it to the constant function 1:

$$\begin{aligned} ([\Delta, A_F] - (2\partial_r^* F'(r)\partial_r + 2r^{-3}F(r)\Delta_Y))1 &= \Delta_g A_F 1 \\ &= -\frac{1}{2}\Delta_g \partial_r^*(F(r)) \end{aligned}$$

We note some particular instances of this computation. Let $\vartheta(r)$ be a function on \mathbb{R} such that

$$\vartheta(r) = \begin{cases} 0 & r \leq 1/2 \\ 1 & r \geq 4, \end{cases}$$

and

$$\vartheta' \geq 0, \quad \vartheta'(r) \geq 1/4 \text{ for } 1 \leq r \leq 2.$$

Let $\Lambda \gg 0$ and \varkappa be a small constant. We compute the following as zero'th order terms for these various cases:

$$\begin{aligned} F(r) = r^s &\implies -\frac{1}{2}\Delta_g \partial_r^*(F(r)) = \frac{n-1}{2}(s-1)(n+s-3)r^{s-3} + O(r^{s-4}), \\ F(r) = (1 - 1/\log r) &\implies -\frac{1}{2}\Delta_g \partial_r^*(F(r)) = (1 + O(1/\log r)) \frac{(n-1)(n-3)}{2r^3} \\ F(r) = \varkappa \vartheta(r/\Lambda) &\implies -\frac{1}{2}\Delta_g \partial_r^*(F(r)) = \varkappa O(r^{-3}) \end{aligned}$$

(with the “big-Oh” term in the last case being uniformly bounded as $\Lambda \rightarrow \infty$).

Thus in particular, we may let $F(r) = r^s$ and (abusing notation) denote A_F by A_s in this case. Then

$$(3.4) \quad [\Delta_g + V, A_s] = 2s\partial_r^* r^{s-1}\partial_r + 2r^{s-3}\Delta_Y - \frac{n-1}{2}(s-1)(n+s-3)r^{s-3} + r^{s-4}Q_Y + f,$$

where Q_Y denotes an operator of second order involving only derivatives in Y with coefficients in $S^0(X)$ and $f \in S^{s-3-\nu}(X)$. The two principal terms have the same (non-negative!) sign if $s \geq 0$, with the sign being definite if $s > 0$; if $n \geq 3$, the zero'th order term has the same sign if in addition $s \leq 1$; the sign is definite if $s < 1$, and in case $n = 3$, $s > 0$.

On the other hand, in dealing with radial derivatives we will employ A_F with $F(r) = (1 - 1/\log r)$. We let A' denote the operator in this case, and compute

$$A' = (1/2) \left(\left(1 - \frac{1}{\log r}\right)A + A \left(1 - \frac{1}{\log r}\right) \right),$$

where

$$A \equiv A_0 = \frac{1}{2}(\partial_r - \partial_r^*).$$

Then we have

$$(3.5) \quad [\Delta_g + V, A'] = 2\partial_r^* \frac{1}{r(\log r)^2} \partial_r + (1 + e_1)2r^{-3}\Delta_Y + (1 + e_2) \frac{(n-1)(n-3)}{2r^3} + r^{-4}Q_Y,$$

with $e_i = O((\log r)^{-1})$ a symbol in $S^0(X)$. (We have recycled the notation Q_Y to denote different operators of the same form.)

Finally, in obtaining ℓ^∞ - ℓ^1 bounds, we will employ

$$F(r) = 1 - \frac{1}{\log r} + \varkappa \vartheta\left(\frac{r}{\Lambda}\right);$$

with \varkappa a (small) positive constant. Let A'' denote the corresponding operator. Thus

$$(3.6) \quad \begin{aligned} [\Delta_g + V, A''] &= 2\partial_r^* \left(\frac{\varkappa}{\Lambda} \vartheta'(r/\Lambda) + \frac{1}{r(\log r)^2} \right) \partial_r \\ &+ (1 + e_1)2r^{-3}\Delta_Y + (1 + e_2) \frac{(n-1)(n-3)}{2r^3} + r^{-4}Q_Y, \end{aligned}$$

with $e_1 = O((\log r)^{-1})$ a symbol in $S^0(X)$ and $e_2 = O((\log r)^{-1}) + \varkappa O(1)$ a symbol in $S^0(X)$, bounded by $1/2$ when $r \gg 0$ and $\varkappa \ll 1$.

We now return to arbitrary metric $g = g_0 + g_1$, with g_0 of the warped product form (3.1). We write Δ_{g_0} for the Laplacian of g_0 and $B^{*,0}$ for the adjoint of an operator B with respect to g_0 , while B^* is written for the adjoint with respect to g . Thus, in (3.2)–(3.5), all adjoints are of the *,0 type, and all Laplacians are that of g_0 . Since A_s and A' depend on the metric via the adjoints, we write $A_{s,0}$ and A'_0 for the g_0 versions. Now, as

$$(3.7) \quad g_1 \in S^{-\nu}(X; {}^{\text{sc}}T^*X \otimes {}^{\text{sc}}T^*X) = S^{2-\nu}(X; {}^{\text{b}}T^*X \otimes {}^{\text{b}}T^*X),$$

we have

$$\sqrt{g} = \sqrt{g_0}(1 + \tilde{g}), \quad \tilde{g} \in S^{-\nu}(X),$$

and thus

$$\partial_r^* - \partial_r^{*,0} \in S^{-1-\nu}(X),$$

i.e. the effect on (3.2) is that e is replaced by a slightly less decaying symbol. Correspondingly,

$$A_s - A_{s,0} \in S^{s-1-\nu} \text{Diff}_b^1(X), \quad A' - A'_0 \in S^{-1-\nu} \text{Diff}_b^1(X), \quad A'' - A''_0 \in S^{-1-\nu} \text{Diff}_b^1(X).$$

Also,

$$\Delta_g - \Delta_{g_0} \in S^{-2-\nu} \text{Diff}_b^2(X).$$

Combining these, and expanding the commutators,

$$(3.8) \quad \begin{aligned} [\Delta_g + V, A_s] - [\Delta_{g_0} + V, A_{s,0}] &\in S^{-3+s-\nu} \text{Diff}_b^2(X), \\ [\Delta_g + V, A'] - [\Delta_{g_0} + V, A'_0] &\in S^{-3-\nu} \text{Diff}_b^2(X) \\ [\Delta_g + V, A''] - [\Delta_{g_0} + V, A''_0] &\in S^{-3-\nu} \text{Diff}_b^2(X), \end{aligned}$$

as always with uniform estimates as $\Lambda \rightarrow \infty$ in the A'' case.

In order to make the foregoing into a global computation, we need to add a cutoff that localizes near infinity. So let $\chi_0(r)$ be a cutoff function equal to 0 for $r < 1$ and 1 for $r > 2$. Let $\chi(r) = \chi_0(r/R)$ with $R \gg 0$. Then

$$(3.9) \quad \begin{aligned} [\Delta_g + V, \chi(r)A_s] &= 2s\partial_r^* \chi(r)r^{s-1}\partial_r + 2r^{s-3}\chi(r)\Delta_Y + \frac{n-1}{2}(1-s)(n+s-3)r^{s-3} \\ &\quad + \nabla_X^*(1-\chi(r))\nabla_X + r^{s-4}E_s + r^{s-3-\nu}\tilde{E}_s, \end{aligned}$$

where $E_s \in \text{Diff}_{\text{sc}}^2(X)$, $\tilde{E}_s \in \text{Diff}_b^2(X)$, and

$$(3.10) \quad \begin{aligned} [\Delta_g + V, \chi(r)A'] &= 2\partial_r^* \frac{\chi(r)}{r(\log r)^2} \partial_r + (1+e_3)2r^{-3}\chi(r)\Delta_Y + (1+e_4)\frac{(n-1)(n-3)}{2r^3} \\ &\quad + \nabla_X^*(1-\chi(r))\nabla_X + r^{-4}E_0 + r^{-3-\nu}\tilde{E}_0 \end{aligned}$$

where $E_0 \in \text{Diff}_{\text{sc}}^2(X)$, $\tilde{E}_0 \in \text{Diff}_b^2(X)$ (and are different from the E_s, \tilde{E}_s in (3.9)), and $e_i = O((\log r)^{-1})$ satisfy symbol estimates. Since on any compact set, the terms with e_i can be absorbed into E_s , we may assume here that $|e_i| < 1/2$ so that

$$1 + e_i > 1/2.$$

For the same reason, by shifting a large compactly supported part of the \tilde{E}_s term in (3.9) into the faster-decaying E_s we may assume that \tilde{E}_s can be absorbed into the first three terms of (3.9) without compromising their positivity: for any $\delta > 0$ we may assume

$$(3.11) \quad \begin{aligned} &|\langle r^{s-3-\nu}\tilde{E}_s v, v \rangle| \\ &\leq \delta \left(2s \|\sqrt{\chi(r)}r^{(s-1)/2}\partial_r v\|^2 + 2\|r^{(s-3)/2}\sqrt{\chi(r)}\nabla_Y v\|^2 \right. \\ &\quad \left. + \frac{n-1}{2}(1-s)(n+s-3)\|r^{(s-3)/2}v\|^2 \right) \end{aligned}$$

for all $v \in \dot{C}^\infty(X)$ (hence by density whenever the right hand side is finite), with. We will take $\delta = 1/2$. Likewise, we may assume that the terms in (3.10) satisfy:

$$(3.12) \quad \begin{aligned} &|\langle r^{-3-\nu}\tilde{E}_0 v, v \rangle| \\ &\leq \delta \left(2\|\frac{\sqrt{\chi(r)}}{r^{1/2}(\log r)}\partial_r v\|^2 + 2\|r^{-3/2}\sqrt{\chi(r)}\nabla_Y v\|^2 + \frac{(n-1)(n-3)}{2}\|r^{-3/2}v\|^2 \right). \end{aligned}$$

We will take $\delta = 1/4$ here.

Exactly the same computation applies in the case of A'' , but with the addition of an extra ∂_r^2 term:

$$(3.13) \quad \begin{aligned} [\Delta_g + V, \chi(r)A''] &= 2\partial_r^* \left(\frac{\chi(r)}{r(\log r)^2} + \frac{\varkappa}{\Lambda} \vartheta'(r/\Lambda) \right) \partial_r + (1 + e_3) 2r^{-3} \chi(r) \Delta_Y \\ &+ (1 + e_4) \frac{(n-1)(n-3)}{2r^3} \\ &+ \nabla_X^* (1 - \chi(r)) \nabla_X + r^{-4} E_0 + r^{-3-\nu} \tilde{E}_0, \end{aligned}$$

with error terms estimated as above (provided \varkappa is taken sufficiently small).

We now consider the pairing

$$(3.14) \quad \int_0^T \langle \chi(r)A'(\square + V)u, u \rangle dt.$$

While we will manipulate this expression for an arbitrary solution u , we remark that the calculations in the following paragraph are unchanged if we replace u by

$$\Psi_H u \equiv \psi(H^2(\Delta_g + V))u,$$

which indeed we will do below.

We first remark that moving $\chi(r)$ and A' to the right slot of the pairing (3.14) and inserting the weight $r^{1/2}(\log r)$, resp. its reciprocal, in the two slots, we obtain, for $\gamma > 0$,

$$(3.15) \quad \begin{aligned} &\left| \int_0^T \langle \chi(r)A'(\square + V)u, u \rangle dt \right| \\ &\leq \left(\int_0^T \|r^{1/2}(\log r)(\square + V)u\|^2 dt \right)^{1/2} \left(\int_0^T \|r^{-1/2}(\log r)^{-1}(A')^* \chi(r)u\|^2 dt \right)^{1/2} \\ &\leq \gamma^{-1} \int_0^T \|r^{1/2}(\log r)(\square + V)u\|^2 dt + \gamma \int_0^T \|r^{-1/2}(\log r)^{-1}(A')^* \chi(r)u\|^2 dt, \end{aligned}$$

where $\langle \cdot, \cdot \rangle$ denotes the spatial inner product. On the other hand, integrating (3.14) twice by parts yields

$$(3.16) \quad -\langle \chi(r)A'u_t, u \rangle|_0^T + \langle \chi(r)A'u, u_t \rangle|_0^T - \int_0^T \langle [\chi(r)A', \Delta_g + V]u, u \rangle dt.$$

Let

$$\mathcal{E}(s) = (\|\nabla_g u\|^2 + \|\partial_t u\|^2 + \|\sqrt{V}u\|^2)|_{t=s}$$

denote the energy norm at fixed time. By the Hardy/Poincaré inequality, we have

$$\|\chi A'u(t)\|^2 \lesssim \mathcal{E}(t).$$

We further compute, as usual, that

$$\mathcal{E}'(t) = 2 \operatorname{Re} \langle u_t, (\square + V)u \rangle \leq 2\mathcal{E}(t)^{1/2} \|(\square + V)u\|.$$

Hence

$$\begin{aligned} \mathcal{E}(T)^{1/2} &\leq \mathcal{E}(0)^{1/2} + \int_0^T \|(\square + V)u\| dt \\ &\lesssim \mathcal{E}(0)^{1/2} + \int_0^T \|(\square + V)u\| dt \end{aligned}$$

Consequently, the first two terms in (3.16) are bounded by a multiple of

$$\mathcal{E}(0) + \left(\int_0^T \|(\square + V)u\| dt \right)^2$$

uniformly in t .³

Thus, (3.10), (3.12) and (3.15) now allow us to estimate (absorbing lower-order terms into positive ones, and choosing γ small so that the second term on the right hand side of (3.15) can be absorbed into the left hand side below):

$$\begin{aligned} (3.17) \quad &\int_0^T \langle \partial_r^* r^{-1} \chi(r) \log(r)^{-2} \partial_r u, u \rangle + \langle r^{-3} u, u \rangle \\ &+ \langle r^{-3} \chi(r) \Delta_Y u, u \rangle + \langle \nabla_X^* \chi(r) \nabla_X u, u \rangle dt \\ &\lesssim \int_0^T \langle r^{-4} E_0 u, u \rangle dt + \mathcal{E}(0) + \left(\int_0^T \|(\square + V)u\| dt \right)^2. \end{aligned}$$

We apply this estimate to spectrally localized data, i.e. replace u by $\Psi_H u$, as indicated at the beginning of the previous paragraph. We state the following lemma more generally than is immediately necessary, in the $n \geq 3$ setting.

Lemma 3.1. *Suppose $n \geq 3$. If $E \in \text{Diff}_{\text{sc}}^2(X)$ and u is as above, then for $0 \leq \rho$, $0 \leq s$, $0 \leq \rho + s < \min(2, n/2)$,*

$$\left| \langle r^{-2(s+\rho)} E \Psi_H u, \Psi_H u \rangle \right| \lesssim H^{-2\rho} \|r^{-s} \Psi_H u\|^2.$$

Proof. It suffices, by rearranging the LHS (as commutator terms require the same form of estimates) to show that for $L \in \text{Diff}_{\text{sc}}^1(X)$,

$$\left\| r^{-(s+\rho)} L \Psi_H u \right\| \lesssim H^{-\rho} \|r^{-s} \Psi_H u\|.$$

Let $\phi \in \mathcal{C}_c^\infty(\mathbb{R})$ be identically 1 on $\text{supp } \psi$. By Proposition 2.1,

$$H^\rho r^{-(s+\rho)} L \phi(H^2(\Delta_g + V)) r^s$$

is uniformly bounded in H , so for all v ,

$$\left\| H^\rho r^{-(s+\rho)} L \phi(H^2(\Delta + V)) v \right\| \lesssim \|r^{-s} v\|.$$

Applying this with $v = \Psi_H u$ completes the proof of the lemma. \square

Taking $s = 3/2$, $\rho = (1 - \epsilon)/2$, $0 < \epsilon < 1$, which we may if $n \geq 4$, we deduce:

³If preferred, we could of course replace the norm on the inhomogeneity by the weighted L^2 spacetime norm

$$\int_0^T \left\| (2 + t^2)^{1/4} \log(2 + t^2) (\square + V)u \right\|^2 dt.$$

Lemma 3.2. *If $n \geq 4$, $E \in \text{Diff}_{\text{sc}}^2(X)$ and u is as above, then for $0 < \epsilon < 1$,*

$$|\langle r^{-4+\epsilon} E \Psi_H u, \Psi_H u \rangle| \lesssim H^{-1+\epsilon} \left\| r^{-3/2} \Psi_H u \right\|^2.$$

As a consequence, we can estimate, a fortiori,

$$\langle r^{-4} E_0 \Psi_H u, \Psi_H u \rangle \lesssim H^{-1+\epsilon} \left\| r^{-3/2} \Psi_H u \right\|^2$$

so that for H sufficiently large, this term may be absorbed in the main term on the left of (3.17). Thus, for H sufficiently large,

$$\begin{aligned} & \int_0^T \left\| r^{-1/2} (\log r)^{-1} \sqrt{\chi}(r) \partial_r \Psi_H u \right\|^2 + \left\| r^{-3/2} \Psi_H u \right\|^2 \\ & \quad + \left\| r^{-3/2} \sqrt{\chi}(r) \nabla_Y \Psi_H u \right\|^2 + \left\| \sqrt{1-\chi} \nabla_X \Psi_H u \right\|^2 dt \\ & \lesssim \mathcal{E}(0) + \left(\int_0^T \|(\square + V)u\| dt \right)^2, \end{aligned}$$

with constants independent of T . This concludes the proof of Theorem 1.1. \square

To prove Theorem 1.1' we proceed similarly, but with A'' replacing A' . We now let

$$(3.18) \quad \Upsilon_k = \{r \in [2^k, 2^{k+1}]\}$$

denote a dyadic decomposition in radial shells. Then we obtain by the same argument

$$\begin{aligned} (3.19) \quad & \left| \int_0^T \langle \chi(r) A'' (\square + V) u, u \rangle dt \right| \\ & = \left| \int_0^T \langle r^{1/2} (\square + V) u, r^{-1/2} (A'')^* \chi(r) u \rangle dt \right| \\ & \leq \int_0^T \sum_k \int_{\Upsilon_k} \left| r^{1/2} (\square + V) u \right| \cdot \left| r^{-1/2} (A'')^* \chi(r) u \right| dg dt \\ & \leq \sum_k \left\| r^{1/2} (\square + V) u \right\|_{L^2([0, T] \times \Upsilon_k)} \left\| r^{-1/2} (A'')^* \chi(r) u \right\|_{L^2([0, T] \times \Upsilon_k)} \\ & \leq \left\| r^{1/2} (\square + V) u \right\|_{\ell^1(\mathbb{N}; L^2([0, T] \times \Upsilon_k))} \left\| r^{-1/2} (A'')^* \chi(r) u \right\|_{\ell^\infty(\mathbb{N}; L^2([0, T] \times \Upsilon_k))} \\ & \leq \gamma^{-1} \left\| r^{1/2} (\square + V) u \right\|_{\ell^1(\mathbb{N}; L^2([0, T] \times \Upsilon_k))}^2 + \gamma \left\| r^{-1/2} (A'')^* \chi(r) u \right\|_{\ell^\infty(\mathbb{N}; L^2([0, T] \times \Upsilon_k))}^2 \end{aligned}$$

Thus, for $\gamma > 0$ we have

$$\begin{aligned}
& \int_0^T \left(\langle \partial_r^* r^{-1} \chi(r) \log(r)^{-2} \partial_r u, u \rangle + \langle \partial_r^* \frac{\mathcal{Z}}{\Lambda} \vartheta'(r/\Lambda) \chi(r) \partial_r u, u \rangle + \langle r^{-3} u, u \rangle \right. \\
& \quad \left. + \langle r^{-3} \chi(r) \Delta_Y u, u \rangle + \langle \nabla_X^* \chi(r) \nabla_X u, u \rangle \right) dt \\
(3.20) \quad & \lesssim \int_0^T \langle r^{-4} E_0 u, u \rangle dt + \mathcal{E}(0) + \left(\int_0^T \|(\square + V)u\| dt \right)^2 \\
& \quad + \gamma^{-1} \left\| r^{1/2} (\square + V)u \right\|_{\ell^1(\mathbb{N}; L^2([0, T] \times \Upsilon_k))}^2 \\
& \quad + \gamma \left\| r^{-1/2} (A'')^* \chi(r) u \right\|_{\ell^\infty(\mathbb{N}; L^2([0, T] \times \Upsilon_k))}^2.
\end{aligned}$$

Now we take the supremum of the LHS over $\Lambda \in \{2^k, k \in \mathbb{N}\}$ to obtain

$$\begin{aligned}
& \int_0^T \left(\langle \partial_r^* r^{-1} \chi(r) \log(r)^{-2} \partial_r u, u \rangle + \langle r^{-3} u, u \rangle + \langle r^{-3} \chi(r) \Delta_Y u, u \rangle \right. \\
& \quad \left. + \langle \nabla_X^* \chi(r) \nabla_X u, u \rangle \right) dt \\
& \quad + \left\| \chi(r) r^{-1/2} \partial_r u \right\|_{\ell^\infty(\mathbb{N}; L^2([0, T] \times \Upsilon_k))}^2 \\
(3.21) \quad & \lesssim \int_0^T \langle r^{-4} E_0 u, u \rangle dt + \mathcal{E}(0) + \left(\int_0^T \|(\square + V)u\| dt \right)^2 \\
& \quad + \gamma^{-1} \left\| r^{1/2} (\square + V)u \right\|_{\ell^1(\mathbb{N}; L^2([0, T] \times \Upsilon_k))}^2 \\
& \quad + \gamma \left\| r^{-1/2} (A'')^* \chi(r) u \right\|_{\ell^\infty(\mathbb{N}; L^2([0, T] \times \Upsilon_k))}^2.
\end{aligned}$$

where we have of course used the fact that $r \sim \Lambda$ on $\text{supp } \vartheta'(r/\Lambda)$. Taking γ sufficiently small to absorb the last term in the LHS yields and applying the resulting estimate to spectrally localized data as above yields Theorem 1.1'. \square

4. A LOCAL MORAWETZ ESTIMATE

In this section, we prove Theorem 1.2. We fix some $\tilde{\psi} \in \mathcal{C}_c^\infty(\mathbb{R})$ which is identically 1 on $\text{supp } \psi$. We write u simply in place of

$$\Psi_H u = \psi(H^2(\Delta + V))u$$

throughout this section to simplify the notation.

As our estimates will be in spacetime, we also define the manifold with corners given by its compactification,

$$M = \overline{\mathbb{R}} \times X,$$

with $\overline{\mathbb{R}}$ denoting the compactification of \mathbb{R} to an interval and X our scattering manifold, regarded as usual as a manifold with boundary endowed with the singular metric (2.1). In this section we will employ the notation $\|\bullet\|_M$ to denote the spacetime L^2 norm, and for the sake of emphasis, we will use $\|\bullet\|_X$ for the spatial norm, previously denoted simply $\|\bullet\|$. We will likewise let ∇_X denote the gradient in the spatial directions only, previous denoted ∇_g , while letting ∇ denote the spacetime gradient.

We may assume that the δ in (1.5) is sufficiently small for convenience (as the result is stronger then); e.g. take $\delta < 1/4$. Let $\phi_0 \in \mathcal{C}_c^\infty(\mathbb{R})$, supported in $[-2, 2]$, identically 1 in $[-1, 1]$. With $0 < c < 3\delta < 1$ to be fixed, let $\phi(\mu) = \phi_0(\mu/c)$, and let $\tilde{\phi}(\mu) = \phi(\mu/(3\delta))$. Also, we let χ be as in Section 3, see the definition just above (3.9), so $\chi \equiv 1$ near infinity. Now, with

$$k = 2\kappa - 1, \quad s = 1 - 2\sigma,$$

let

$$B_{k,s} = \frac{1}{2} \left(\phi(r/t)^2 \tilde{\phi}(t)^2 (r/t)^s t^{-k} \chi(r) \partial_r - \partial_r^* \phi(r/t)^2 \tilde{\phi}(t)^2 (r/t)^s t^{-k} \chi(r) \right),$$

so on the complement of $\text{supp}(1 - \phi\tilde{\phi})$ (where $\phi\tilde{\phi} \equiv 1$) we actually have

$$B_{k,s} = t^{-k-s} A_s$$

with A_s given by (3.3). Note that $\text{supp } d\tilde{\phi} \cap \text{supp } \phi$ is a compact subset of M° , so we can effectively ignore terms in which $\tilde{\phi}$ is commuted through differential operators below.

Before proceeding, we recall that in $\{r > r_0, t > 1\} \subset M$, $\mathcal{V}_b(M)$ is spanned by $r\partial_r, t\partial_t$ and vector fields on Y , over $\mathcal{C}^\infty(M)$. We now blow up the corner $\partial_2 M$, where both $x = 0$ and $t^{-1} = 0$, to obtain the manifold

$$\tilde{M} = [M; \partial_2 M].$$

The blowup procedure (cf. [11, Appendix]; the simple setting of [10, Section 4.1] where a codimension 2 corner is blown up is analogous to the present case, with our t^{-1} taking the place of x' in [10]) serves to replace the corner $t^{-1} = r^{-1} = 0$ of M by its inward-pointing spherical normal bundle, in this case diffeomorphic to Y times an interval. This new boundary hypersurface will be denoted⁴ *mf*. We let *tf* denote the ‘‘temporal infinity’’ face given by the lift of the set $t^{-1} = 0$ in M to \tilde{M} . The main consequence of the blowup procedure is that it introduces r/t as a smooth function where it is bounded, and its differential is non-vanishing on the interior of *mf* (where it is thus one of the standard coordinates, namely the variable along the interval referred to above; others being t^{-1} and y , coordinates on $Y = \partial X$; we use t^{-1} as a boundary defining function for *mf*). However, $\mathcal{V}_b(\tilde{M})$ is still spanned by the lift of the same vector fields, except that the smooth structure is replaced by $\mathcal{C}^\infty(\tilde{M})$. Correspondingly, the symbol spaces are unaffected by the blow-up. Now, the support of $d\phi$ only intersects the front face of \tilde{M} among all boundary faces, and r/t is bounded from both above and below by positive constants there. Thus, on the support of $d\phi$, $t^{-1}\mathcal{V}_b(\tilde{M})$ is locally generated by ∂_r, ∂_t and $r^{-1}\partial_Y$, i.e. by vector fields corresponding to the energy space. Thus, on $\text{supp } d\phi$, $\nabla v \in t^\kappa L^2(M)$ implies $v \in t^{\kappa+1} L^2(M)$, $k < (n-1)/2$, by a Hardy inequality, and so we obtain more generally:

⁴The notation is short for ‘‘main face,’’ as for X Euclidean (diffeomorphically, not metrically), this can be identified with an open dense subset of the boundary of the radial compactification of Minkowski space. Note that the light cone, $r/t = 1$, hits the boundary in the interior of *mf*. If $r/t = 1$ in *mf* were blown up inside \tilde{M} , the resulting front face is where Friedlander’s radiation field [4] can be defined by rescaling a solution to the wave equation. Indeed, in the interior of this front face (in $t > 0$) $r - t = (r/t - 1)/t^{-1}$ becomes a smooth variable along the fibers of the front face.

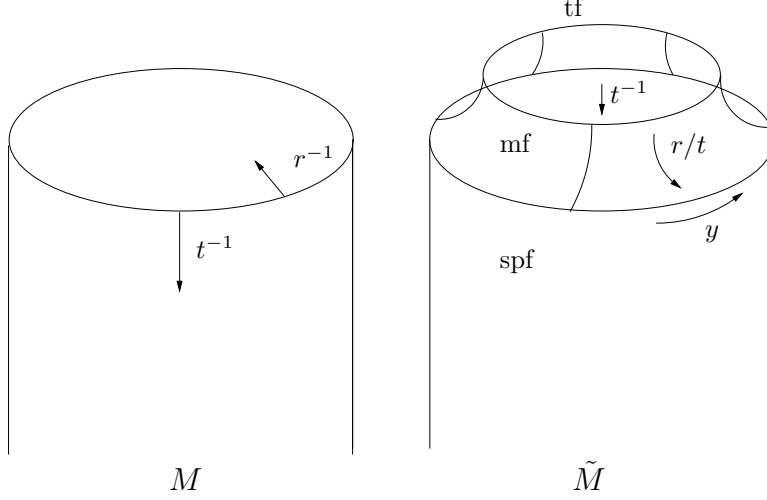


FIGURE 1. The compactified spacetime M and the blown-up spacetime \tilde{M} .

Lemma 4.1. *Let Υ be an open set in \tilde{M} with $\bar{\Upsilon} \cap \partial\tilde{M} \subset \text{mf}^\circ$. If $\text{supp } v \subset \Upsilon$,*

$$\nabla v \in t^\kappa L^2(M) \implies v \in t^{\kappa+1} H_b^1(\tilde{M}).$$

Here the Sobolev space is still relative to the metric density (so $L^2(M)$ and $L^2(\tilde{M})$ are the same, as only the smooth structure at infinity changed).

Before proceeding, we caution the reader that *near temporal infinity*, tf , i.e. the lift of $t = +\infty$ on M to \tilde{M} , weighted versions of $\mathcal{V}_b(\tilde{M})$ do not give rise to the finite length vector fields relative to $dt^2 + g$ (i.e. the energy space). Indeed, in the interior of tf , one is in a compact set of the spatial slice X , and $\mathcal{V}_b(\tilde{M})$ is spanned by $t\partial_t$ and smooth vector fields on X , while the energy space corresponds to ∂_t and smooth vector fields on X . In order to emphasize the structure when we stay away from tf and also from spatial infinity spf , i.e. the lift of $\partial X \times \bar{\mathbb{R}}$, we write

$$\tilde{M}' = \tilde{M}^\circ \cup \text{mf}^\circ = \tilde{M} \setminus (\text{tf} \cup \text{spf}).$$

Our convention below is that inner products and norms are on M unless otherwise specified by a subscript X (in which case they are on X). The inner product on M on functions is given by integration against the density of $dt^2 - g$, or equivalently that of $dt^2 + g$. The inner product on vectors on M uses the positive definite inner product, based on $dt^2 + g$, unless otherwise specified. The one case where we take advantage of the “otherwise specified” disclaimer is when we consider the indefinite Dirichlet form in (4.4), where we use the indefinite form induced by $dt^2 - g$ (so the density is positive definite, but not the pointwise pairing between vectors) in order to use the PDE.

Now,

$$(4.1) \quad [\partial_t^2, t^{-k-s}] = -2(k+s)t^{-k-s-1}\partial_t + (k+s)(k+s+1)t^{-k-s-2} \in t^{-k-s-2}\text{Diff}_b^1(\bar{\mathbb{R}}_t).$$

Moreover, on \tilde{M}' (hence on $\text{supp } d\phi$),

$$B_{k,s} \in t^{-k-1}\text{Diff}_b^1(\tilde{M}'), \quad \square \in t^{-2}\text{Diff}_b^2(\tilde{M}') \implies [\square, B_{k,s}] \in t^{-k-3}\text{Diff}_b^2(\tilde{M}').$$

In particular, from (3.9),

$$\begin{aligned}
(4.2) \quad & [\partial_t^2 + \Delta_g + V, B_{k,s}] \\
& = t^{-k-s} \phi(r/t)^2 \tilde{\phi}(t)^2 \left(\sum_i Q_i^* r^s t^{-1} R_i \right. \\
& \quad + (2s \partial_r^* r^{s-1} \chi(r) \partial_r + 2r^{s-3} \chi(r) \Delta_Y + \frac{n-1}{2} (1-s)(n+s-3) r^{s-3} \\
& \quad \left. + \nabla_X^* (1-\chi(r)) \nabla_X + r^{s-4} E_s + r^{s-3-\nu} \tilde{E}_s \right) + F_{k,s},
\end{aligned}$$

where $Q_i \in t^{-1} \text{Diff}_b^1(\mathbb{R}_t)$, $R_i \in r^{-1} \text{Diff}_b^1(X)$ contains no y -derivatives and is supported on $\text{supp } \chi$, $E_s \in \text{Diff}_{\text{sc}}^2(X)$, $\tilde{E}_s \in \text{Diff}_b^2(X)$, $F_{k,s} \in t^{-k-3} \text{Diff}_b^2(\tilde{M})$ and $F_{k,s}$ is supported on $\text{supp } d\phi$. Note that the term with ∇_X has spatially compact support, and could be absorbed into E_s , but writing the commutator in the stated manner is convenient for it gives a positive definite result modulo E_s , which we later control separately, by low energy techniques. Moreover, for $0 < s < 1$ (which is the case if $0 < \sigma < 1/2$), we may assume that the \tilde{E}_s term can be absorbed into the first three terms on the right hand side by (3.11). Thus, for $0 < s < 1$ the spatial term, arising from (3.9), is positive modulo the E_s term, while the terms $Q_i^* r^s t^{-1} R_i$ arising from the commutator with ∂_t^2 , are indefinite. However, we can estimate

$$\begin{aligned}
& \|t^{-(k+s)/2} r^{(s-1)/2} \tilde{\phi} \phi R_i u\|_M \\
& \leq C_0 (\|t^{-(k+s)/2} r^{(s-1)/2} \tilde{\phi} \phi \partial_r u\|_M + \|t^{-(k+s)/2} \tilde{\phi} \phi r^{(s-3)/2} u\|_M),
\end{aligned}$$

with C_0 independent of c in the definition of ϕ . We also have a similar estimate for the time derivatives, using the PDE and the fact that $t^{-1} \lesssim r^{-1}$ on $\text{supp } \phi$:

$$\begin{aligned}
(4.3) \quad & \|t^{-(k+s)/2} r^{(s-1)/2} \tilde{\phi} \phi Q_i u\|_M \\
& \leq C_1 (\|t^{-(k+s)/2} r^{(s-1)/2} \tilde{\phi} \phi \partial_t u\|_M + \|t^{-(k+s)/2} r^{(s-3)/2} \tilde{\phi} \phi u\|_M) \\
& \leq C_2 \left(\|t^{-(k+s)/2} r^{(s-1)/2} \tilde{\phi} \phi \nabla_X u\|_M + \|t^{-(k+s)/2} r^{(s-3)/2} \tilde{\phi} \phi u\|_M \right. \\
& \quad + \|t^{-(k+s)/2} r^{(s+1)/2} \tilde{\phi} \phi (\square + V) u\|_M \\
& \quad \left. + \|t^{-(k+s)/2} r^{(s-1)/2} \nabla_X u\|_{\text{supp } d(\phi\tilde{\phi})} + \|t^{-(k+s)/2} r^{(s-3)/2} \tilde{\phi} \phi u\|_{\text{supp } d(\phi\tilde{\phi})} \right).
\end{aligned}$$

To show the validity of the last step, consider the (indefinite!) Dirichlet form relative to $dt^2 - g$, and use that

$$\begin{aligned}
(4.4) \quad & \langle t^\alpha r^\beta \tilde{\phi} \phi \nabla u, t^\alpha r^\beta \tilde{\phi} \phi \nabla u \rangle_{M, dt^2 - g} - \langle t^\alpha r^{\beta+1} \tilde{\phi} \phi \square u, t^\alpha r^{\beta-1} \tilde{\phi} \phi u \rangle_M \\
& = \langle f_0 t^{\alpha-1} r^\beta \tilde{\phi} \phi u, t^\alpha r^\beta \tilde{\phi} \phi \partial_t u \rangle_M + \langle f t^\alpha r^{\beta-1} \tilde{\phi} \phi u, t^\alpha r^\beta \tilde{\phi} \phi \nabla_X u \rangle_M \\
& \quad + \sum_j \langle t^{\alpha-1} r^\beta Q'_j u, t^\alpha r^\beta R'_j u \rangle_M,
\end{aligned}$$

where f_0 , f and Q'_j are zero'th order symbols on \tilde{M} , $R'_j \in t^{-1} \mathcal{V}_b(\tilde{M})$, and Q'_j is supported on $\text{supp } d(\phi\tilde{\phi})$ (so R'_j can be taken supported nearby). Indeed, the terms on the right hand side can be controlled: the second term directly by the positive spatial term via Cauchy-Schwarz,

$$(4.5) \quad |\langle f t^\alpha r^{\beta-1} \tilde{\phi} \phi u, t^\alpha r^\beta \tilde{\phi} \phi \nabla_X u \rangle_M| \lesssim \|t^\alpha r^{\beta-1} \tilde{\phi} \phi u\|_M^2 + \|t^\alpha r^\beta \tilde{\phi} \phi \nabla_X u\|_M^2,$$

while the first can be Cauchy-Schwarz with a small constant in front of the ∂_t factor in the pairing, which then can be reabsorbed in the Dirichlet form, while the other factor can be controlled directly from the spatial positive term using that r/t is bounded, and indeed small:

$$(4.6) \quad |\langle f_0 t^{\alpha-1} r^\beta \tilde{\phi} \phi u, t^\alpha r^\beta \tilde{\phi} \phi \partial_t u \rangle_M| \lesssim \epsilon \|t^\alpha r^\beta \tilde{\phi} \phi \partial_t u\|_M^2 + \epsilon^{-1} \|t^{\alpha-1} r^\beta \tilde{\phi} \phi u\|_M^2.$$

In addition, the $\square u$ term can be controlled using Cauchy-Schwarz:

$$(4.7) \quad \langle t^\alpha r^{\beta+1} \tilde{\phi} \phi \square u, t^\alpha r^{\beta-1} \tilde{\phi} \phi u \rangle_M \leq \|t^\alpha r^{\beta-1} \tilde{\phi} \phi u\|_M^2 + \|t^\alpha r^{\beta+1} \tilde{\phi} \phi \square u\|_M^2.$$

Finally, we change from \square to $\square + V$ using that $V \in S^{-2}(X)$, so

$$(4.8) \quad \|t^\alpha r^{\beta+1} \tilde{\phi} \phi V u\|_M \lesssim \|t^\alpha r^{\beta-1} \tilde{\phi} \phi u\|_M.$$

Combining (4.4)–(4.8),

$$(4.9) \quad \begin{aligned} & (1 - \tilde{C}\epsilon) \|t^\alpha r^\beta \tilde{\phi} \phi \partial_t u\|_M^2 \\ & \lesssim \|t^\alpha r^\beta \tilde{\phi} \phi \nabla_X u\|_M^2 + \epsilon^{-1} \|t^{\alpha-1} r^\beta \tilde{\phi} \phi u\|_M^2 + (\|t^\alpha r^{\beta-1} \tilde{\phi} \phi u\|_M^2 + \|t^\alpha r^\beta \tilde{\phi} \phi \nabla_X u\|_M^2) \\ & \quad + (\|t^{\alpha-1} r^\beta u\|_{\text{supp } d(\phi\tilde{\phi})}^2 + \|t^\alpha r^\beta \tilde{\phi} \phi \nabla_X u\|_{\text{supp } d(\phi\tilde{\phi})}^2) \\ & \quad + \|t^\alpha r^{\beta+1} \tilde{\phi} \phi (\square + V) u\|_M^2, \end{aligned}$$

where $\epsilon > 0$ can be arbitrarily chosen without affecting \tilde{C} , so in particular we may assume $\tilde{C}\epsilon < 1/2$. This proves the second inequality in (4.3).

Since

$$(4.10) \quad \begin{aligned} & |\langle Q_i^* \tilde{\phi} \phi r^s t^{-1} t^{-(k+s)} R_i u, u \rangle_M| \\ & \leq 2c \|t^{-(k+s)/2} r^{(s-1)/2} \tilde{\phi} \phi R_i u\|_M \|t^{-(k+s)/2} r^{(s-1)/2} \tilde{\phi} \phi Q_i u\|_M \\ & \leq c \|t^{-(k+s)/2} r^{(s-1)/2} \tilde{\phi} \phi R_i u\|_M^2 + c \|t^{-(k+s)/2} r^{(s-1)/2} \tilde{\phi} \phi Q_i u\|_M^2, \end{aligned}$$

with c arising from the extra factor of r/t , which is $\leq 2c$ on $\text{supp } \phi$, in the pairing over what is included in the two factors on the right, for sufficiently small $c > 0$ these can be absorbed in the positive definite spatial term in (4.2), modulo terms supported on $\text{supp } d(\phi\tilde{\phi})$ and modulo terms involving $(\square + V)u$.

It remains to check that the ‘‘quantum’’ term (in that we need global estimates on tf to control it), E_s , in (4.2) can be controlled. With $\rho_0 \in C_c^\infty(\mathbb{R})$ identically 1 on a neighborhood of $\text{supp } \phi_0$, supported in $(-3, 3)$, and let $\rho(\mu) = \rho_0(\mu/c)$, c as above, and consider $\rho = \rho(r/t)$. We write

$$|\langle r^{s-4} \phi(r/t) E_s \phi(r/t) u, u \rangle_X| \lesssim \|r^{(s-4)/2} \phi(r/t) \nabla_X u\|_X^2 + \|r^{(s-4)/2} \phi(r/t) u\|_X^2.$$

We remark that the left hand side rearranges the factors from (4.2) by commuting a factor of ϕ through E_s ; the difference is supported in $\text{supp } d\phi$, and when multiplied by t^{-k-s} , it is in $t^{-k-5} \text{Diff}_b^2(\tilde{M})$, i.e. can be controlled the same way $F_{k,s}$ was (and is indeed better behaved). Now write

$$u = \tilde{\psi}(H^2(\Delta + V))\rho(r/t)u + \tilde{\psi}(H^2(\Delta + V))(1 - \rho(r/t))u,$$

and estimate the summands (and their X -gradients) individually.

First, to deal with the $\tilde{\psi}(H^2(\Delta + V))(1 - \rho(r/t))u$ term, note that by Proposition 2.9, for any N ,

$$\begin{aligned} & \|\phi(r/t) \tilde{\psi}(H^2(\Delta + V))(1 - \rho(r/t))r^N\|_{\mathcal{L}(L^2(X), L^2(X))} \\ & + \|\nabla_X \phi(r/t) \tilde{\psi}(H^2(\Delta + V))(1 - \rho(r/t))r^N\|_{\mathcal{L}(L^2(X), L^2(X))} \leq C_{H,N} t^{-N}, \end{aligned}$$

where the constant $C_{H,N}$ depends on H (and N). Since u is tempered, for sufficiently large N , the space-time norm of $r^{-N}t^{-N}u$ is bounded (with the bound depending on H), controlling this term.

Next, to deal with the

$$\tilde{u} = \tilde{\psi}(H^2(\Delta + V))\rho(r/t)u$$

term, we use the low energy localization, i.e. we take H large. First, for $\epsilon' \in (0, 1/2)$,

$$\begin{aligned} \|r^{(s-4)/2}\phi(r/t)\tilde{u}\|_X^2 &\leq \|r^{(s-4)/2}\tilde{u}\|_X^2 \leq \|r^{(s-3)/2-\epsilon'}\tilde{u}\|_X^2, \\ \|r^{(s-4)/2}\phi(r/t)\nabla_X\tilde{u}\|_X^2 &\leq \|r^{(s-3)/2-\epsilon'}\nabla_X\tilde{u}\|_X^2. \end{aligned}$$

By the Hardy/Poincaré lemma, Lemma 3.1, using $0 < s < 1$, and choosing $\epsilon' \in (0, 1/2)$ such that $-(s-3)/2 + \epsilon' < 3/2$,

$$\begin{aligned} &\|r^{(s-3)/2-\epsilon'}\tilde{u}\|_X^2 + \|r^{(s-3)/2-\epsilon'}\nabla_X\tilde{u}\|_X^2 \\ (4.11) \quad &\lesssim H^{-2\epsilon'}\|r^{(s-3)/2}\rho(r/t)u\|_X^2 \\ &= H^{-2\epsilon'}(\|\phi(r/t)r^{(s-3)/2}u\|_X^2 + \|(\rho(r/t) - \phi(r/t))r^{(s-3)/2}u\|_X^2). \end{aligned}$$

For H sufficiently large, the first term on the right hand side of (4.11) can be absorbed into the zero'th order positive definite spatial term,

$$(4.12) \quad \left\langle \phi(r/t)^2(2s\partial_r^*r^{s-1}\chi(r)\partial_r + 2r^{s-3}\chi(r)\Delta_Y + \frac{n-1}{2}(1-s)(n+s-3)r^{s-3})u, u \right\rangle_X,$$

modulo errors supported on $\text{supp } d(\tilde{\phi}\phi)$, which are of the same kind as those in $F_{k,s}$ in (4.2). On the other hand, the second term on the right hand side of (4.11) is supported on $\text{supp}(\rho - \phi)$, and is thus similar to the $F_{k,s}$ terms, with slightly larger support (but still within the region where this proposition gives no improvement upon the a priori assumption). Explicitly, when also integrated in t , on $\text{supp}(\rho - \phi)$ we have $\nabla u \in t^\kappa L^2(M)$ (with $(k+1)/2 = \kappa$), and since on $\text{supp}(\rho - \phi)$, $r \sim t$,

$$\begin{aligned} &\int t^{-k-s}\tilde{\phi}(t)^2\|(\rho(r/t) - \phi(r/t))r^{(s-3)/2}u\|_X^2 dt \\ (4.13) \quad &\lesssim \int t^{-k-s}\tilde{\phi}(t)^2\|t^{(s-3)/2}(\rho(r/t) - \phi(r/t))u\|_X^2 dt \\ &\lesssim \|t^{-(k+1)/2}\tilde{\phi}(t)^2(\rho(r/t) - \phi(r/t))u\|_M^2. \end{aligned}$$

Now, for $\gamma > 0$,

$$\begin{aligned} &|\langle B_{k,s}u, (\square + V)u \rangle_M| = |\langle B_{k,s}u, \rho(r/t)(\square + V)u \rangle_M| \\ &\leq \|r^{(s+1)/2}t^{-(k+s)/2}\rho(r/t)(\square + V)u\|_M \|r^{-(s+1)/2}t^{(k+s)/2}B_{k,s}u\|_M \\ &\leq \gamma^{-1}\|r^{(s+1)/2}t^{-(k+s)/2}\rho(r/t)(\square + V)u\|_M^2 + \gamma\|r^{-(s+1)/2}t^{(k+s)/2}B_{k,s}u\|_M^2 \\ &\leq \gamma^{-1}\|r^{(s+1)/2}t^{-(k+s)/2}\rho(r/t)(\square + V)u\|_M^2 + \gamma\left(\|r^{(s-1)/2}t^{-(k+s)/2}\rho(r/t)\chi(r)\partial_r u\|_M^2 \right. \\ &\quad \left. + \|r^{(s-3)/2}t^{-(k+s)/2}\rho(r/t)u\|_M^2 + \|t^{-(k+s)/2}(1-\chi(r))\rho(r/t)\nabla_X u\|_M^2\right). \end{aligned}$$

Taking $\gamma > 0$ sufficiently small, the second term on the right hand side can be absorbed into the positive spatial terms in (4.2), modulo terms that are supported

outside $\text{supp}(1 - \phi\tilde{\phi})$, but in that region we have a priori control of these terms. We thus deduce from (4.2) that

$$(4.14) \quad \begin{aligned} & \|t^{-(k+s)/2}r^{(s-1)/2}\tilde{\phi}\phi\nabla_X u\|_M + \|t^{-(k+s)/2}r^{(s-3)/2}\tilde{\phi}\phi u\|_M \\ & \lesssim \|r^{(s+1)/2}t^{-(k+s)/2}(\square + V)u\|_{L^2(\Omega)} + \|u\|_{t^{(k+3)/2}H_b^1(\Upsilon)}, \end{aligned}$$

where $\Upsilon = \text{supp } d\phi \cup \text{supp}(\rho - \phi)$ and with the last space being $H_b^1(\tilde{M})$ on Υ , provided $0 < s < 1$ and provided the right hand side is finite. Recall that Lemma 4.1 means that the second term on the right hand side may be replaced by

$$\|\nabla u\|_{t^{(k+1)/2}L^2(\Upsilon)}.$$

In view of the Dirichlet form estimate, (4.9), we also get the bound for the time derivative which is analogous to (4.14). Noting that

$$\begin{aligned} & \|t^{-(k+s)/2}r^{(s-1)/2}\tilde{\phi}\phi\nabla u\|_M + \|t^{-(k+s)/2}r^{(s-3)/2}\tilde{\phi}\phi u\|_M \\ & = \|t^{-(k+1)/2}(t/r)^{(1-s)/2}\tilde{\phi}\phi\nabla u\|_M + \|t^{-(k+1)/2}(t/r)^{(1-s)/2}r^{-1}\tilde{\phi}\phi u\|_M, \end{aligned}$$

and $0 < s < 1$, so $0 < (1-s)/2 < 1/2$, while (using $r/t < 1$ on Ω and $s > 0$)

$$\|r^{(s+1)/2}t^{-(k+s)/2}(\square + V)u\|_{L^2(\Omega)} \lesssim \|r^{1/2}t^{-k/2}(\square + V)u\|_{L^2(\Omega)},$$

this proves the estimate of the theorem, provided the formal pairings employed above are finite and hence the computation justified.

In order to employ the above argument without a priori assumptions on u , we need to introduce a weight, $(1 + \alpha t)^{-1}$, $\alpha > 0$, in $B_{k,s}$ to justify the arguments, and let $\alpha \rightarrow 0$. Roughly speaking, such a weight does not cause problems since we could deal with *arbitrary* weights t^{-k} already (i.e. k did not need to have a sign). More precisely, the contribution of this weight to the commutator of $(1 + \alpha t)^{-1}B_{k,s}$ and \square is similar to (4.1), namely $[\partial_t^2, t^{-k-s}(1 + \alpha t)^{-1}] \in t^{-k-s-3}\text{Diff}_b^1(\mathbb{R}_t)$ for $\alpha > 0$, and it is uniformly bounded in the larger space $t^{-k-s-2}\text{Diff}_b^1(\mathbb{R}_t)$ for $\alpha \in [0, 1)$. Thus, its role is analogous to that of the $Q_i^*R_i$ terms of (4.2), which in turn arose from (4.1); these are estimated in (4.10), and can be controlled by making c small. Letting $\alpha \rightarrow 0$, and using standard functional analytic arguments, completes the proof of the theorem. \square

We finally remark that a version of this argument also works for solutions of the wave equation at intermediate frequencies, via Mourre estimate techniques (see especially [5]). While this is not interesting for the homogeneous wave equation, whose middle-frequency solutions will have rapid decay inside the light cone, for the inhomogeneous equation it may be of some interest. The key point then is that one works with $\psi(\Delta + V)$, and $\tilde{\psi}(\Delta + V)$, supported near a fixed energy $\lambda_0 > 0$, and one controls the analogue of the E_s term (that arose above) by shrinking the support of $\tilde{\psi}$, using that $\tilde{\psi}(\Delta + V)E_s$ converges to 0 in norm as the support is shrunk to $\{\lambda_0\}$ by the relative compactness of E_s (which follows from elliptic estimates and its decay at spatial infinity). This holds since $\tilde{\psi}(\Delta + V)$ converges to 0 strongly as the support is shrunk to $\{\lambda_0\}$, since λ_0 is not an L^2 eigenvalue of $\Delta + V$.

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